C\(^0\)- Continuity Isoparametric Formulation using Trigonometric Displacement Functions for One Dimensional Elements

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C0- Continuity Isoparametric Formulation using Trigonometric Displacement Functions for One Dimensional Elements

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Abstract- This is an original research on the selection of the trigonometric shape functions in the finite element analysis of the one dimensional elements. A new family of C0- continuity elements is introduced using the trigonometric interpolation model. To relate the natural and global coordinate system for each element of specific structure (i.e. transformation mapping) in one dimensional element a new trigonometric function is used and the new determinant has been introduced instead of polynomial function and Jacobian determinant. The new introduced trigonometric determinant allows for the state of constant strain within the element satisfying the completeness requirement. However, this cannot be achieved using the Jacobian determinant to relate the coordinates while using the trigonometric functions. The finite element formulation presented in this paper gives comparable results with exact solution for all kinds of one dimensional analysis.

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1. Introduction

Finite element method (FEM) is the approximate piecewise analysis in the domain of interest, researchers have put in efforts to select an appropriate interpolating function which can very closely approximate the field variable and converge to the exact solution. Polynomials have been studied for many years, starting in the 19th century, and they have shown to have mostly good approximation properties. Nevertheless, they are not “good for all seasons” [1]. In [2], it was shown that for differential equations with rough coefficients, the finite element method using polynomial shape functions can lead to arbitrarily “bad” results. Effective shape functions should have good approximation properties in entire domain of the interest. To increase the accuracy of the solution various procedures for error estimation have been devised and mesh refinement is used. Various procedures exist for the refinement of finite element (FE) solutions. More researches have been reported on the references [4-14].

By considering the linear-strain triangular (LST) element it can be seen that the development of element matrices and equations expressed in terms of a global coordinate system becomes an enormously difficult task [15]. The isoparametric method may appear somewhat tedious (and confusing initially), but it leads to a simple computer program formulation, and it is generally applicable for one-, two- and three-dimensional stress analysis and for nonstructural problems. The isoparametric formulation allows elements to be created that are nonrectangular and have curved sides [16].

In this paper, we first illustrate the trigonometric isoparametric formulation to develop the shape functions of C0 continuity of the family of one dimensional bar elements and to derive the strain matrix, stiffness matrix and then force vector. Use of the bar element makes it relatively easy to understand the method because it involves simple expressions. Then quantitative concepts for assessing and comparing effectiveness of these families of shape functions are given. Focus on the principles that should govern the selection of the trigonometric shape functions are discussed, and one dimensional elements are studied by employing these new shape functions obtained from trigonometric displacement functions to analyze the bars carrying the self-weight and the results have been compared with the exact solutions of classical methods of solid mechanics.

II. Isoparametric Concept and Coordinate Systems

The term isoparametric is derived from the use of the same shape functions (or interpolation functions) to define the element’s geometric shape as are used to define the displacements within the element. Isoparametric element equations are formulated using a natural (or intrinsic) coordinate system T that is defined by element geometry and not by the element orientation in the global coordinate system. In other words, axial coordinate T is attached to the bar and remains directed along the axial length of the bar, regardless of how the bar is oriented in space [16]. The relationship between the natural coordinate system T and the global coordinate system X for each element of specific structure is called the transformation mapping and must be used in the element equation formulations. The coordinate systems are shown in fig. 1.
The natural coordinate system $T$ is a dimensionless quantity varying from $T_1$ to $T_2$ at node 1 and node 2 respectively. In natural coordinate system the position of any point inside the element is varying by $\sin(\alpha T)$. The natural coordinate system is attached to the element, with the origin located at its center, as shown in Fig. 1(b). The $T$ axis needs not be parallel to the $x$ axis, this is only for convenience.

For the special case consider a circle of unit radius shown in Fig. 2, when the $T$ and $x$ axes are inside the circle and parallel to each other. The $T$ and $x$ axes having the origin located at the center of the element are coincided at the center of the circle ($X_c = \frac{x_1 + x_2}{2}$). For the special case when $\alpha = \frac{\pi}{2}$ and the $-1 \leq T \leq 1$ and $-1 \leq x \leq 1$ the global and natural coordinates can be related by

$$X = X_c + \frac{L}{2} \sin(\frac{\pi}{2} T)$$  \hspace{1cm} (1)

Where $X_c$ is the global coordinate of the element's centroid.

The displacement function within the bar which relates the displacement at any point inside the element to the nodal displacements is given by

$$U = \sum N U_i$$ \hspace{1cm} (2)

The function which relates the coordinate of any point within the element to the global coordinate is given by

$$X = \sum N X_i$$ \hspace{1cm} (3)

By using the equation (3) the shape functions have been used for coordinate transformation from natural coordinate system to the global Cartesian system and mapping of the parent element to required shape in global system successfully achieved. This formula is given by Taig [17].

In Eq. (3) the summation is over the number of nodes of the element. $N$ is the shape function, $U_i$ are the nodal displacements and $X_i$ is the coordinates of nodal points of the element. The shape functions are to be expressed in natural coordinate system.

The equations (2) and (3) can be written in matrix form as

$$\{U\} = [N] \{U\}_e$$ \hspace{1cm} (4)

$$\{X\} = [N] \{X\}_e$$ \hspace{1cm} (5)

Where $\{U\}$ is vector of displacement at any point, $\{U\}_e$ is vector of nodal displacements, $\{X\}_e$ is the vector of nodal coordinates and $\{X\}$ is the vector of coordinate of any point in global system.

### III. Interpolation Model and Shape Functions for Two Nodded Element

The quality of approximation achieved by Rayleigh–Ritz and FE approaches depends on the admissible assumed trial, field or shape functions. These functions can be chosen in many different ways. The most universally preferred method is the use of simple polynomials. It is also possible to use other functions such as trigonometric functions [18, 19]. While choosing the interpolation model and shape functions, the following considerations have to be taken into account[3, 20].

a) To ensure convergence to the correct result certain simple requirements must be satisfied as following criteria.

**Criterion 1.** The displacement shape functions chosen should be such that they do not permit straining of an
element to occur when the nodal displacements are caused by a rigid body motion.

**Criterion 2.** The displacement shape functions have to be of such forms that if nodal displacements are compatible with a constant strain condition such constant strain will in fact be obtained.

**Criterion 3.** The displacement shape functions should be chosen such that the strains at the interface between elements are finite (even though they may be discontinuous).

b) The pattern of variation of the field variable resulting from the interpolation model should be independent of the local coordinate system.

c) The number of generalized coordinates should be equal to the number of nodal degrees of freedom of the element.

The interpolation model of the field variable (displacement model inside the element) in terms of nodal degrees of freedom is given by trigonometric sine function as

\[ U(T) = a_1 + a_n \sin\left(\frac{\pi}{2} T\right) \quad \text{Where} \quad -1 \leq T \leq 1 \]  

Where \( a_1 \) and \( a_n \) are the coefficients known as generalized coordinates and must be equal to the number of nodal unknowns \( M \). In equation (6), the nodal values of the solution, also known as nodal degrees of freedom, are treated as unknowns in formulating the system or overall equations. To express the interpolation model in terms of the nodal degrees of freedom of a typical finite element \( e \) having \( M \) nodes, the values of the field variable at the nodes can be evaluated by substituting the nodal coordinates into the Eq. (6). The Eq. (6) can be expressed in general form of

\[
\begin{bmatrix}
U(\text{at node 1}) \\
U(\text{at node 2})
\end{bmatrix}^{(e)} = \begin{bmatrix}
\bar{U}(1) \\
\bar{U}(2)
\end{bmatrix}^{(e)} = \begin{bmatrix}
1 & \sin\left(-\frac{\pi}{2}ight) \\
1 & \sin\left(\frac{\pi}{2}\right)
\end{bmatrix} \begin{bmatrix}
a_1 \\
a_2
\end{bmatrix} = \begin{bmatrix}
1 & -1 \\
1 & 1
\end{bmatrix} \begin{bmatrix}
a_1 \\
a_2
\end{bmatrix} = [\eta] \bar{\alpha}
\]

Where \( \bar{U}^{(e)} \) is the vector of nodal values of the field variable corresponding to element \( e \), and the square matrix \( [\eta] \) can be identified from Eq. (9). By inverting Eq. (9), we obtain

\[
\begin{bmatrix}
\bar{U}_1 \\
\bar{U}_2
\end{bmatrix} = \begin{bmatrix}
1 & \sin\left(-\frac{\pi}{2}\right) \\
1 & \sin\left(\frac{\pi}{2}\right)
\end{bmatrix}^{-1} \begin{bmatrix}
a_1 \\
a_2
\end{bmatrix} = [\eta]^{-1} \alpha
\]
\[ \vec{a} = \left[ \eta \right]^{-1} \vec{\eta}^{(e)} \quad (10) \]

Substituting the Eq. (10) Into Eq. (7) gives
\[ \vec{U} = \vec{\eta}^T \vec{a} = \vec{\eta}^T \left[ \eta \right]^{-1} \vec{\eta}^{(e)} = [N]\vec{\eta}^{(e)} \quad (11) \]

Thus \([N] = \vec{\eta}^T \left[ \eta \right]^{-1} \quad (12)\]

Where, \([N]\) is the matrix of interpolation functions or shape functions.

Equation (11) expresses the interpolating function inside any finite element in terms of the nodal unknowns of that element, \(\vec{\eta}^{(e)}\). A major limitation of trigonometric interpolation functions is that one has to invert the matrix \([\eta]\) to find \(\vec{U}\), and \([\eta]^{-1}\) may become singular in some cases.

\[ N_1 = \frac{\sin \left( \frac{\pi}{2} T \right) - \sin \left( \frac{\pi}{2} \right)}{\sin \left( \frac{\pi}{2} \right) - \sin \left( -\frac{\pi}{2} \right)} \]
\[ N_2 = \frac{\sin \left( \frac{\pi}{2} \right) - \sin \left( \frac{\pi}{2} \right)}{\sin \left( \frac{\pi}{2} \right) - \sin \left( -\frac{\pi}{2} \right)} \quad (13) \]

Therefore, the shape functions are
\[ \begin{cases} N_1 = \frac{1 - \sin \left( \frac{\pi}{2} T \right)}{2} \\ N_2 = \frac{\sin \left( \frac{\pi}{2} \right) + 1}{2} \end{cases} \quad (14) \]

It must be noted that \(-1 \leq T \leq 1\).

The variation of the resulting shape functions are shown in Fig. 4. The essential properties of shape functions are that they must be unity at one node and zero at the other nodes. It can be seen that by shifting the \(T\) to \(T_1\) and \(T_2\) we get
\[ \begin{cases} At \ node \ 1 \ where \ T = T_1 = -1 \quad \text{At node 2 where } \ T = T_2 = 1 \\ N_1 = 1 \quad N_1 = 0 \\ N_2 = 0 \quad N_2 = 1 \end{cases} \]

Therefore, the natural coordinate system it can be written as
\[ U = N_1 \vec{U}_1 + N_2 \vec{U}_2 \quad (15) \]

To have the \(C^0\) continuity element the sum of the shape functions must be 1 (i.e. \(\sum N_i = 1\)) and the first derivative of the field variable must be zero (i.e. \(\sum \frac{\partial N_i}{\partial T} = 0\)). As there are two nodal unknowns \(U_1\) and \(U_2\) for node 1 and node 2 respectively, therefore in the natural coordinate system it can be written as
\[ \begin{cases} \frac{\partial N_1}{\partial T} = \frac{\partial N_1}{\partial T} + \frac{\partial N_2}{\partial T} = 0 \\ \frac{\partial N_1}{\partial T} = \frac{-\pi}{2} \cos \left( \frac{\pi}{2} T \right) + \frac{\pi}{2} \cos \left( \frac{\pi}{2} T \right) \end{cases} \]

And
\[ \begin{cases} \frac{\partial N_1}{\partial T} = \frac{\partial N_1}{\partial T} + \frac{\partial N_2}{\partial T} = 0 \\ \frac{\partial N_1}{\partial T} = \frac{-\pi}{2} \cos \left( \frac{\pi}{2} T \right) + \frac{\pi}{2} \cos \left( \frac{\pi}{2} T \right) \end{cases} \]
It can be seen that the two essential requirements of the C0 continuity element are satisfied.

It is of interest to mention that there is clear difference between the interpolation model of the element $\mathbf{U}_0 = \mathbf{N} \mathbf{d}$ that applies to the entire element and expresses the variation of the field variable inside the element in terms of the generalized coordinates $\mathbf{a}_i$ and the shape function $\mathbf{N}_i$ that corresponds to the $i^{th}$ nodal degree of freedom $\mathbf{U}_i^e$ and only the sum $\sum \mathbf{N}_i \mathbf{U}_i^e$ represents the variation of the field variable inside the element in terms of the nodal degrees of freedom $\mathbf{U}_i^e$. In fact, the shape function corresponding to the $i^{th}$ nodal degree of freedom $\mathbf{N}_i$ assumes a value of 1 at node $i$ and 0 at all the other nodes of the element[20].

b) Mapping of the element in global coordinate system

The mapping of the parent element in global coordinate system can be done by using eq. (2) which can be written in matrix form as

$$\{\mathbf{X}\} = \{\mathbf{N}_1 \ \mathbf{N}_2\} \begin{bmatrix} \chi_1 \\ \chi_2 \end{bmatrix}$$

(16)

It is clear that $\frac{d\mathbf{u}}{dT} = \left[ \begin{array}{c} -\frac{\pi}{2} \cos \left(\frac{\pi}{2} T\right) \\ \frac{\pi}{2} \cos \left(\frac{\pi}{2} T\right) \end{array} \right] \begin{bmatrix} U_1 \\ U_2 \end{bmatrix}$ and $\frac{d\chi}{dT} = \frac{L}{2} \pi \cos \left(\frac{\pi}{2} T\right)$, therefore the Eq. (17) becomes

$$\mathbf{\epsilon} = \frac{d\mathbf{u}}{dx} = \frac{1}{L} \left[ \begin{array}{c} -1 \\ 1 \end{array} \right] \begin{bmatrix} U_1 \\ U_2 \end{bmatrix}$$

(18)

Strain displacement relation is given as [3]

$$\mathbf{\epsilon} = \sum \mathbf{B}_i \mathbf{\epsilon}_i$$

Or in matrix form as

$$\{\mathbf{\epsilon}\} = \{\mathbf{B}\}_{i=1} \mathbf{\epsilon}_i$$

(20)

Where, $\{\mathbf{\epsilon}\}$ is strain at any point in the element, $\{\mathbf{U}\}_{i}$ is displacement vector of nodal values of the element and $[\mathbf{B}]$ is strain displacement matrix.

By comparing the Eq. (20) with expression given for the strain in Eq. (18) we have the strain displacement matrix of the bar as

$$[\mathbf{B}] = \frac{1}{L} \left[ \begin{array}{cc} -1 & 1 \end{array} \right]$$

(21)

The essential necessity of linear interpolation functions is that the strain must be constant inside the element for with C1- continuity. As it can be seen in Eq. (21) the strain is constant and is same as the strain matrix for bar element using the polynomial functions.

The stress strain relation is given by constitutive

$$\{\sigma\} = [\mathbf{D}] \{\mathbf{\epsilon}\} = [\mathbf{D}] \{\mathbf{B}\}_{i=1} \mathbf{\epsilon}_i$$

(22)

$$\frac{dx}{dT} = \frac{L}{2} \pi \cos \left(\frac{\pi}{2} T\right) \therefore \frac{L}{2} \pi \cos \left(\frac{\pi}{2} T\right) dT = dx$$

(25)
By inserting Eq. (25) in Eq. (24), we can write the stiffness matrix in global coordinate system as

$$[k] = \int_0^L [B]^T [D] [B] A \frac{L \pi}{2} \cos \left(\frac{\pi}{2}T\right) dT$$

(26)

Or

$$[k] = \int_0^L [B]^T [D] [B] |J| dT$$

(27)

Where, $|J| = \frac{dx}{dT} = \frac{\pi L}{2} \cos \left(\frac{\pi}{2}T\right)$ is the Jacobian determinant for one dimensional element with trigonometric displacement functions and relates an element length in the global coordinate system to an element length in the natural coordinate system. This is different from the Jacobian determinant for one dimensional element with polynomial displacement function given by $\frac{L}{2}$ but the concept is same.

By inserting the modulus of elasticity $E= [D]$, Eq. (27) becomes

$$[k] = \int_0^L [B]^T E [B] A \frac{L \pi}{2} \cos \left(\frac{\pi}{2}T\right) dT$$

(28)

By substituting the strain displacement matrix given in Eq. (21), the stiffness matrix can be evaluated as

$$[k] = \frac{EA}{L} \left[ \begin{array}{cc} 1 & -1 \\ -1 & 1 \end{array} \right]$$

(29)

It can be realized that the stiffness matrix is the same as that of given for the two nodded bar element evaluated employing the polynomial functions.

d) Derivation of the system equations in terms of the natural coordinate system

The body and surface forces in terms of the natural coordinate system can be evaluated by the following formulas

$$(F)^e = \iiint_{V} [N]^T [X_b] dV - \iiint_{S} [N]^T [T_s] dS$$

(30)

Where, the product of $A$ and $L$ represents the volume of the element and $X_b$ the body force per unit volume, then $[X_b] A \frac{L \pi}{4} \cos \left(\frac{\pi}{2}T\right)$ is the total body force acting on the element and $T_s$ is traction force-per-unit-length, $[T_s] A \frac{L \pi}{4} \cos \left(\frac{\pi}{2}T\right)$ is now the total traction force. The element equilibrium equation is

$$[K] (U)^e = (F)^e$$

(33)

The above equation of equilibrium is to be assembled for entire structure and boundary conditions are to be introduced. Then the solutions of equilibrium equations result into nodal displacements of all the nodal points. Once these basic unknowns are found, then displacement at any point may be obtained by Eq. (11), the strains may be assembled using the Eq. (12) and then stresses also can be found using the Eq. (22).
The variation of the resulting shape functions are shown in Fig. 5.

Figure 5: Variation of shape functions for bar element in natural and global coordinate system.

To relate the natural coordinate \( T \) (where, \( 0 \leq T \leq 1 \)) and global coordinate \( X \) (where, \( 0 \leq x \leq 1 \)) the Jacobian determinant given in Eq. (25) becomes

\[
\frac{dx}{dT} = \frac{ln}{2} \cos(\frac{\pi}{2} T) \quad \text{therefore} \quad \frac{ln}{2} \cos(\frac{\pi}{2} T) \, dT = dx
\]  

(36)

The strain displacement matrix \( [B] \) will be same as given in Eq. (21) and the stiffness matrix \( [K] \) same as Eq. (29). The consistent forces will be as

\[
\{F\}^s = \iint_x \{N\}^T \{X\} A \frac{ln}{2} \cos(\frac{\pi}{2} T) \, dT - \iint_x \{N\}^T \{T_x\} \frac{ln}{2} \cos(\frac{\pi}{2} T) \, dT
\]  

(37)

It must be noted that the limits of the integrations will be 0 to 1.

IV. INTERPOLATION MODEL AND SHAPE FUNCTIONS FOR THREE NODDED ELEMENT

To illustrate the concept of three nodded elements using the trigonometric functions, the element with three coordinates of nodes, \( x_1, x_2, \) and \( x_3, \) in the global coordinates is shown in Fig. 6. Again the element is considered within a circle of unit radius and the third node is selected at the centre of the circle.

The interpolation model of the field variable (displacement model inside the element) in terms of nodal degrees of freedom is given by trigonometric function as

\[
U(T) = a_1 + a_2 \sin \left(\frac{\pi}{2} T\right) + a_3 \left(\sin \left(\frac{\pi}{2} T\right)\right)^2 \quad \text{Where} \quad \begin{cases} T = -1 & \text{at node one} \\ T = 1 & \text{at node two} \\ T = 0 & \text{at node three} \end{cases}
\]  

(38)

Using the displacement field given in Eq. (38), the shape functions are given as

\[
\begin{align*}
N_1 &= \frac{(\sin \left(\frac{\pi}{2} T\right))^2 - \sin^2 \left(\frac{\pi}{2} T\right)}{2} \\
N_2 &= \frac{(\sin \left(\frac{\pi}{2} T\right))^2 + \sin^2 \left(\frac{\pi}{2} T\right)}{2} \\
N_3 &= 1 - (\sin \left(\frac{\pi}{2} T\right))^2
\end{align*}
\]  

(39)
The variation of the resulting shape functions are shown in Fig. 7. The essential properties of shape functions are also satisfied as following.

\[
\begin{align*}
\text{At node 1 where } T = T_1 = -1 \\
N_1 &= 1 \\
N_2 &= 0 \\
N_3 &= 0
\end{align*}
\begin{align*}
\text{At node 2 where } T = T_2 = 1 \\
N_1 &= 0 \\
N_2 &= 1 \\
N_3 &= 0
\end{align*}
\begin{align*}
\text{At node 3 where } T = T_3 = 0 \\
N_1 &= 0 \\
N_2 &= 0 \\
N_3 &= 1
\end{align*}
\]

\[
U = N_1 \bar{U}_1 + N_2 \bar{U}_2 + N_3 \bar{U}_3
\]

\[
\sum N_i = 1
\]

\[
\sum \frac{\partial N_i}{\partial T} = 0
\]

\[
\frac{\partial N_1}{\partial T} = \frac{\pi}{2} \cos\left(\frac{\pi}{2} T\right) \sin\left(\frac{\pi}{2} T\right) - \frac{\pi}{4} \cos\left(\frac{\pi}{2} T\right) + \frac{\pi}{2} \cos\left(\frac{\pi}{2} T\right) \sin\left(\frac{\pi}{2} T\right) + \frac{\pi}{4} \cos\left(\frac{\pi}{2} T\right) - \pi \cos\left(\frac{\pi}{2} T\right) \sin\left(\frac{\pi}{2} T\right) = 0
\]

It can be seen that the two essential requirements of the \(C^0\) continuity element are satisfied.

\(a\) Strain – displacement and stress - strain relationship

From our basic definition of axial strain we have

\[
\epsilon = \frac{du}{dx} = \frac{dU}{dT} = [B] \begin{bmatrix} U_1 \\ U_2 \\ U_3 \end{bmatrix}^{(e)}
\]

(41)

It has previously proven that \(\frac{dx}{dT} = \frac{L_0}{4} \cos\left(\frac{\pi}{2} T\right)\), this relationship holds for the three nodded one-dimensional elements as well as for the two-noded...
\[ \frac{du}{dx} = \left[ \frac{4}{L \cos \left( \frac{\pi}{2} T \right)} \left( \frac{\pi}{2} \cos \left( \frac{\pi}{2} x \right) \sin \left( \frac{\pi}{2} x \right) - \frac{\pi}{4} \cos \left( \frac{\pi}{2} T \right), \frac{\pi}{2} \cos \left( \frac{\pi}{2} T \right) \sin \left( \frac{\pi}{2} T \right) + \frac{\pi}{4} \cos \left( \frac{\pi}{2} T \right), -\pi \cos \left( \frac{\pi}{2} T \right) \sin \left( \frac{\pi}{2} T \right) \right] \right] \begin{bmatrix} U_1 \\ U_2 \\ U_3 \end{bmatrix}^{(e)} \]

Therefore

\[ \frac{du}{dx} = \frac{1}{L} \begin{bmatrix} 2 \sin \left( \frac{\pi}{2} T \right) - 1, \ 2 \sin \left( \frac{\pi}{2} T \right) + 1, \ -4 \sin \left( \frac{\pi}{2} T \right) \end{bmatrix} \begin{bmatrix} U_1 \\ U_2 \\ U_3 \end{bmatrix}^{(e)} \] (42)

By comparing the expression given for the strain in Eq. (41) with Eq. (19), the strain-displacement matrix \([B]\) for the three nodded bar is

\[ [B] = \frac{1}{L} \begin{bmatrix} 2 \sin \left( \frac{\pi}{2} T \right) - 1, \ 2 \sin \left( \frac{\pi}{2} T \right) + 1, \ -4 \sin \left( \frac{\pi}{2} T \right) \end{bmatrix} \] (43)

Substituting the expression for \([B]\) into Eq. (27), the stiffness matrix is obtained as

\[ [K] = \frac{1}{L} \begin{bmatrix} \frac{2 \sin \left( \frac{\pi}{2} T \right) - 1}{2 \sin \left( \frac{\pi}{2} T \right) + 1} \left[ 2 \sin \left( \frac{\pi}{2} T \right) - 1, \ 2 \sin \left( \frac{\pi}{2} T \right) + 1, \ -4 \sin \left( \frac{\pi}{2} T \right) \right] \frac{L \pi}{4} \cos \left( \frac{\pi}{2} T \right) dT \]

\[ = \frac{EA}{L} \int_{-1}^{1} \begin{bmatrix} 2 \sin \left( \frac{\pi}{2} T \right) - 1 \\ 2 \sin \left( \frac{\pi}{2} T \right) + 1 \end{bmatrix} \begin{bmatrix} 2 \sin \left( \frac{\pi}{2} T \right) - 1, \ 2 \sin \left( \frac{\pi}{2} T \right) + 1, \ -4 \sin \left( \frac{\pi}{2} T \right) \end{bmatrix} \frac{\pi}{4} \cos \left( \frac{\pi}{2} T \right) dT \]

\[ = \frac{EA}{L} \int_{-1}^{1} \begin{bmatrix} (2 \sin \left( \frac{\pi}{2} T \right) - 1)^2 & (2 \sin \left( \frac{\pi}{2} T \right))^2 - 1 & -8 \sin \left( \frac{\pi}{2} T \right)^2 + 4 \sin \left( \frac{\pi}{2} T \right) \\ (2 \sin \left( \frac{\pi}{2} T \right))^2 - 1 & (2 \sin \left( \frac{\pi}{2} T \right) + 1)^2 & -8 \sin \left( \frac{\pi}{2} T \right)^2 - 4 \sin \left( \frac{\pi}{2} T \right) \\ -8 \sin \left( \frac{\pi}{2} T \right)^2 + 4 \sin \left( \frac{\pi}{2} T \right) & -8 \sin \left( \frac{\pi}{2} T \right)^2 - 4 \sin \left( \frac{\pi}{2} T \right) & (4 \sin \left( \frac{\pi}{2} T \right))^2 \end{bmatrix} \frac{\pi}{4} \cos \left( \frac{\pi}{2} T \right) dT \]

Integrating the matrix the stiffness matrix for three nodded bar element becomes

\[ [K] = \frac{EA}{L} \begin{bmatrix} 2.333333 & 0.333333 & -2.666667 \\ 0.333333 & 2.333333 & -2.666667 \\ -2.666667 & -2.666667 & 5.333333 \end{bmatrix} \] (44)

The stiffness matrix given in Eq. (27) is the same as that of given for the three nodded bar elements evaluated using polynomial functions.

**Example 1.** Analysis of bar of uniform cross section \((A)\), Young’s modulus of the material \((E)\) due to self-weight (unit weight, \(\rho\)) when held as shown in Fig. 8. Self-weight acting in \(T\) direction.

\[
\begin{align*}
\text{Figure 8: Bar of constant cross section}
\end{align*}
\]
\[ \varepsilon = \frac{\rho L}{2k} \]  
\[ \sigma = \frac{\rho L}{2} \]  

Equations (45) and (46) are the exact solutions for a bar having constant cross-section due to its own self weight.

V. Conclusion

Using the trigonometric interpolation model, new family of \( C^0 \)-continuity elements are introduced. To obtain the constant stress and strain state in 2 nodded elements, trigonometric function is used instead of the polynomial Jacobian determinant to relate the natural and global coordinate system. The bar of uniform cross section is analyzed and results are compared with those of obtained using the polynomial functions.

References Références Referencias


