

Computational Analysis of Combustion Chamber Using Cavity-based fuel Injector with Non-Premixed Combustion Model

K.M.Pandey¹ and A.P.Singh²

¹ National Institute of Technology, Silchar, Assam, India.

Received: 10 December 2011 Accepted: 1 January 2012 Published: 15 January 2012

Abstract

This paper presents the supersonic combustion of hydrogen fuel using cavity-based fuel injector with two-dimensional turbulent non-premixed combustion model. The present model is based on the standard k-epsilon (two equations) with standard wall functions which is P1 radiation model and a PDF (Probability Density Function) approach is created. The hydrogen fuel is injected just upstream of the cavity. The Contour of Mass fraction of OH indicates a little amount of OH around 0.001454 after combustion. A cavity flame holder is provided which injects hydrogen fuel in a supersonic hot air stream that facilitates enhanced mixing and combustion efficiency.

Index terms— Mach number, CFD, combustion, hydrogen fuel, non-premixed combustion, scramjet, standard k-epsilon turbulence model, standard wall functions, steady st

1 Introduction

The future of hypersonic air-breathing vehicles lies in the successful development and design of Supersonic combustion ramjet (SCRAMJET) engines which poses some major challenges that has attracted the attention and imagination of researchers worldwide. The serious issues like fuel-air mixing, flame holding, pressure losses and thermal loading can be resolved with the successful implementation of a fuel injection system that provides rapid mixing between the fuel and oxidizer streams, induces pressure losses to a minimum with reduced or zero adverse effects on flame holding capability or thermal/structural integrity of the device. A very short time for fuel injection, fuel-air mixing and subsequently combustion is available of the order of 1 ms and hence the increasing need to develop a system that effectively integrates fuel injection and flame holding for supersonic combustion exists. Thus cavity flame holders has been proposed in recent years as a new concept for flame holding and stabilization in supersonic combustors [4].

Some recent publications have brought to light the subject of cavity flows and their relevance to flame holding in supersonic combustion engines [1,2,3]. Low speed combustion studies with an axis symmetric cavity [5] found optimum flame holding performance using a cavity with its length-to-depth ratio $L=D$ sized for the minimum aerodynamic drag. Longer cavities produced vortex shedding that resulted in unstable flames, and shorter cavities did not provide enough air entrainment to hold the flame. Experimental and numerical results were shown to agree closely on this point [6]. Cavities with small aspect ratios provide better flame holding capability than longer cavities with aft ramp angles as suggested in a study by Yu et al [7] where fuel was injected upstream of a variable L/D cavity at flow speed of Mach 2.

A configuration having a baseline fuel injector/ flame holder with a low angled fuel injection upstream of a wall cavity was used by Tarun Mathur et al [8] where fuel injection and flame piloting was done in a scramjet combustor with all the components contained in the wall. In contrast to in-stream concepts that introduce additional friction drag, wave drag, and cooling requirements to the combustor, this configuration uses no in-stream devices, thereby minimizing these detrimental effects and simplifying the overall combustor and system

designs. Similar studies which involves flush-wall injection upstream of similar cavities in non reacting supersonic flow have provided valuable insights into the effects of cavity configuration (L=D ratio, offset ratio, aft ramp angle), fuel injection pressure, and imposed back pressure on drag, residence times, and fuel distribution within the cavity [9,10]. The combustion experiments as described by Tarun Mathur et al [8] as well as some numerical simulations of cavity-based fuel injector/flame holder [11,12,13] have shown robust flame holding and combustion performance in a scramjet combustor simulating Mach 4-6 conditions at a dynamic pressure of 47.9 k Pa. Some difficulties associated with hydrocarbon fuels which primarily include the relatively long ignition delay time and the challenge in diffusing stable combustion energy into the main flow without disturbing the flow and creating drag penalties may be tackled by cavity-based flame holders as suggested by Ben Yakar et al [2]. A cavity-based flame holder a) creates a sheltered subsonic recirculation area of hot combustion products and increases the effective residence time for the fuel, and b) acts as a pilot light to spread hot combustion products into the main flow. The flow in the vicinity of the cavity can be very stable and can limit the amount of mass entrainment. As can be seen from the fig. 1 below which is a result of numerical computations by Gruber et al [10] there are trapped vortices within the cavity, including a large primary recirculation zone that interacts with the free stream, and a smaller fuel-rich secondary vortex in the forward corner of the cavity. The Cavity flow regimes has been categorized basically into two types by Ben Yakar et al [14] that depends primarily on length-to-depth ratio, L/D. In all the cases it is seen that a shear layer gets separated from the upstream lip and get reattached downstream. The reattachment takes place in the back face for $L/D < 7-10$ and hence are termed as open. For $L/D < 2-3$ transverse oscillation mechanism plays the dominant role but large aspect ratio cavities are controlled by longitudinal oscillations. The high pressure at the rear face as a result of the shear layer impingement increases the drag of the cavity. For $L=D > 10-13$ the cavity flow is termed "closed" because the free shear layer reattaches to the lower wall. The pressure increase in the back wall vicinity and the pressure decrease in the front wall results in large drag losses. The critical length-to-depth ratio, at which a transition between different cavity flow regimes occurs, depends also on the boundary-layer thickness at the leading edge of the cavity, the flow Mach number, and the cavity width.

Another way of improving fuel the fuel-air mixture within the cavity can be direct fuel injection into the cavity as investigated by Allen et al [15]. This resulted in decreased size of fuel rich vortex with subsequent improvement in combustion within the cavity which was due to improved fuel air mixture because of additional air injected directly into the cavity. They also observed that the air injection technique did not have merely a undeviating effect on the fuel-rich region, in fact increasing the air injection without bound had diminishing effect, and eventually are verse effect. For lower fuel injection rates, if the air injection was increased to its maximal limit the combustion increases seen at lower air injection rates moderated to levels near the original fuel-only case. It would seem that the direct air injection technique is able to cause the cavity fuel-air mixture to become too lean to gain any enhancements in combustion if the air injection rate is not organized.

2 II.

3 Materials and Methods

4 a) Physical Model

A mathematical model consists of equations concerning the dependent and the independent variables and the relevant parameters that describe some physical phenomenon. In general, a mathematical prototype consists of differential equations that govern the performance of the physical system, and the related boundary conditions which is shown in figure 2. The advantage of employing the complete Navier-Stokes equations extends not only the investigations that can be carried out on a wide range of flight conditions and geometries, but also in the process the location of shock wave, as well as the physical characteristics of the shock layer, can be exactly determined. We begin by describing the threedimensional forms of the Navier-Stokes equations below. Note that the two-dimensional forms are just simplification of the governing equations in the three dimensions by the omission of the component variables in one of the co-ordinate directions. Neglecting the presence of body forces and volumetric heating, the three-dimensional Navier-Stokes equations are derived as [16]: Continuity Equation: (5) Assuming a Newtonian fluid, the normal stress τ_{xx} , τ_{yy} and τ_{zz} can be taken as combination of the pressure p and the normal viscous stress components τ_{xx} , τ_{yy} , and τ_{zz} while the remaining components are the tangential viscous stress components whereby $\tau_{xy} = \tau_{yx}$, $\tau_{xz} = \tau_{zx}$, and $\tau_{yz} = \tau_{zy}$. For the energy conservation for supersonic flows, the specific energy, E is solved instead of the usual thermal energy H applied in subsonic flow problems. In three dimensions, the specific energy E is repeated below for convenience: $E = e +$

5

It is evident from above that the kinetic energy term contributes greatly to the conservation of energy because of the high velocities that can be attained for flows, where $Ma > 1$. Equations (??)-(??) represent the form of governing equations that are adopted for compressible flows. The solution to the above governing equations nonetheless requires additional equations to close the system. First, the equation of state on the assumption of a perfect gas unemploved, that is,

6 $P = \rho R T$

where R is Gas constant Second, assuming that the air is calorically perfect, the following relation holds for the internal energy: $e = C_v T$

where C_v is specific heat at constant volume. Third, if the Prandtl number is assumed constant (approximately 0.71) for calorically perfect air), the thermal conductivity can be evaluated by the following: The Sutherland's law is typically used to evaluate viscosity μ , which is provided by: $\mu = \mu_0 \left(\frac{T}{T_0} \right)^{1.5} \left(\frac{T_0 + S}{T + S} \right)^2$

Where μ_0 and T_0 are the reference values at standard sea level conditions

7 Computational and Model Parameters a) Geometry and mesh generation

Mesh generation was performed in a Fluent preprocessing program called Gambit. The current model is cavity-based fuel injector with non-premixed combustion as shown in figure 3. The boundary conditions are such that, the air inlet and fuel inlet surfaces are both defined as pressure inlets and the outlet is defined as pressure outlet. Recent research has revealed that perhaps the numerical model will improve if the air inlet is defined as pressure inlet and the fuel inlet is defined as a mass flow inlet. In this particular model the walls of the combustor duct do not have thicknesses. The domain is completely contained by the combustor itself; therefore there is actually no heat transfer through the walls of the combustor. During analysis we have taken same pressure for both fuel and air for all the models. Pressure inlet and pressure outlet conditions were taken on the left and right boundaries respectively. Pressure inlet condition was taken for fuel injector. The top and bottom boundaries, which signify the sidewalls of the isolator, had symmetry conditions on them. The walls, obstacles and other materials were set to standard wall conditions. The computations were initially carried out with various levels of refinement of mesh. There exists a definite level of refinement beyond which there is no significant quantitative change in the result. The limit of that refinement is called the Grid Independent Limit (GIL). The input parameters that were for the model is shown in tabulated form.

8 c) Modeling Details

In the CFD model, the Standard $k-\epsilon$ turbulent model is selected which is one of the most common turbulence models. It is a two equation model that means it includes two extra transport equations to represent the turbulent properties of the flow. This two equation model accounts for history effects like convection and diffusion of turbulent energy. Further, because of the intense turbulent combustion, the eddydissipation reaction model is adopted. The eddydissipation is based on the hypothesis of infinitely fast reactions and the reaction rate is controlled by turbulent mixing. Both the Arrhenius rate and the mixing rate are calculated and the smaller of the two rates is used for the turbulent combustion. While no-slip conditions are applied along the wall, but due to the flow being supersonic, at the outflow all the physical variables are extrapolated from the internal cells. Energy equations were considered and the solution was initialized from the air inlet for simplicity. For hydrogen-air mixing, ideal gas mixing law was followed for determination of thermal conductivity and viscosity, while density was assumed to be for ideal gas. Mass diffusivity was assumed to be following kinetic theory.

IV.

9 Results and Discussions

The various plots of properties such as static temperature, densities etc. along the length of the combustor for the different models are given below. The red colored regions are the regions where the properties attain their maximum values. The blue colored regions indicate the regions where the properties are at their minimum. The properties that were analyzed were: The static temperature was taken as an indication of combustion efficiency of the fuel (hydrogen). Higher combustion efficiency means a greater percentage of the injected fuel undergoes combustion resulting in a higher static temperature at the combustor exit. Study of the mass fraction contours of H_2 , O_2 and H_2O showed evidence of fuel injection, air fuel mixing and combustion respectively. The presence of H_2O indicated the occurrence of combustion. Turbulent kinetic energy was an indication of vortex formation in the cavity which enhances air-fuel mixing. The X-velocity was the velocity at which the combustion products exit the combustor. It represented the thrust available for propulsion of the scramjet. The static pressure and density contours and static pressure and density graphs help in visualizing the shock waves produced by the velocity of hydrogen injection. Moreover, interaction of the reflected shock waves with the air-fuel mixing boundary (visible in the density and static pressure contours) further enhanced the mixing and promoted.

10 a) Static Temperature

From Fig 4 it is evident that static temperature increases from inlet to the outlet. This is due to combustion of the air and injected H_2 fuel. The heat released due to combustion heats up the combustion products (water) and hence, an increase in the static temperature from 398K to 1789 K is observed. The below graph shows the distribution of H_2 in the interior of the combustor. As can be seen, the mass fraction of hydrogen is maximum at the fuel injection port and continues to decrease along the length of the combustor due to combustion. Thus, the graph provides evidence of combustion. Typically, when dealing the chemical reaction, it's important to

11 CONCLUSION

159 remember that mass is conserved, so the mass of product is same as the mass of reactance. Even though the
160 element exists in different the total mass of each chemical element must be same on the both side of equation.
161 The contour and XY Plot of O₂ Mass fraction for the flow field downstream of the injector is shown in the figure
162 12 and figure 13. Oxygen is increased in every combustion reaction in combustion applications and air provides
163 the required oxygen. All components other than air collected together with nitrogen. In air 21% of oxygen and
164 79% of nitrogen are present on a molar basis. From the figure 12 it is observed that, the maximum mass fraction
165 of O₂ is 0.213 which is seen at the beginning of combustion. Figure 13 shows that the profile between the mass
166 fraction of O₂ and the position of the combustion on all conditions such as air inlet, fuel inlet, pressure outlet,
167 default interior and all walls. The contour of mass fraction of OH is shown in figure 14. From the figure 14 it
168 is observed that, the maximum mass fraction of OH is 0.001454 which is found out after combustion, where the
169 minimum value is 0. Figure 15 shows that the profile between the mass fraction of OH and the position of the
170 combustion on all conditions such as air inlet, fuel inlet, pressure outlet, default interior and all walls.

11 Conclusion

171
172 The computational analysis of 2D cavity-based fuel injector was carried out with k- ϵ turbulence model for
173 exposing the flow structure of progress of hydrogen jet through the areas disturbed by the reflections of oblique
174 shock. For that single step reaction kinetics has been used to model the chemistry. The k- ϵ turbulence model
175 also predicted the fluctuations in those regions where the turbulence is reasonably isotropic. From the maximum
176 mass fraction of OH a very small amount of OH (1.45e-03) was observed after combustion. From the above
177 analysis, it is observed that for a scramjet engine having a wall injector with a cavity of L/D=5, if hydrogen is
178 injected at a speed of Mach 1.5 to an incoming air stream at Mach 3.12 speed, a rich air-fuel mixture can be
179 achieved and efficient combustion of this mixture gives a maximum temperature of 1789K at the outlet of the
180 combustor. Also, there is a weak shock formation. Hence, better flame holding can be achieved if the wall injector
181 is coupled with a cavity having with an L/D ratio of 5. Due to ever increasing human need for greater speed
182 and reduced travel time, hypersonic combustion systems will become more and more important in the future. As
183 the mixing time for fuel in the combustor system is very less (~1ms), newer and better injection systems have
184 to be developed that enhance fuel-air mixing and reduce ignition delay period, thus increasing both combustion
efficiency and thrust. ^{1 2 3 4}



1

Figure 1: Fig. 1 :



2

Figure 2: Fig. 2 :

185

¹© 2012 Global Journals Inc. (US) ? ?

²(A) © 2012 Global Journals Inc. (US) Computational Analysis of Combustion Chamber Using Cavity-based fuel Injector with Non-Premixed Combustion Model

³© 2012 Global Journals Inc. (US) Fig 6 : Contour of Density

⁴June Computational Analysis of Combustion Chamber Using Cavity-based fuel Injector with Non-Premixed Combustion Model



3

Figure 3: Fig 3 :



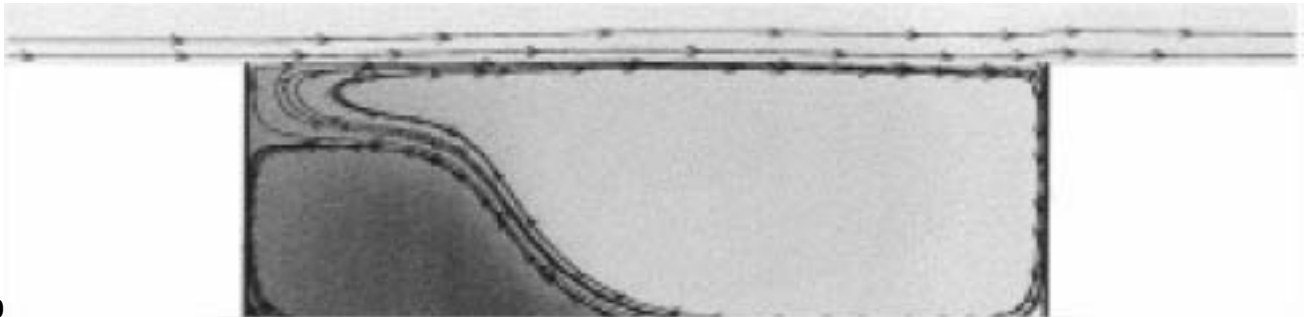
4

Figure 4: Fig 4 :



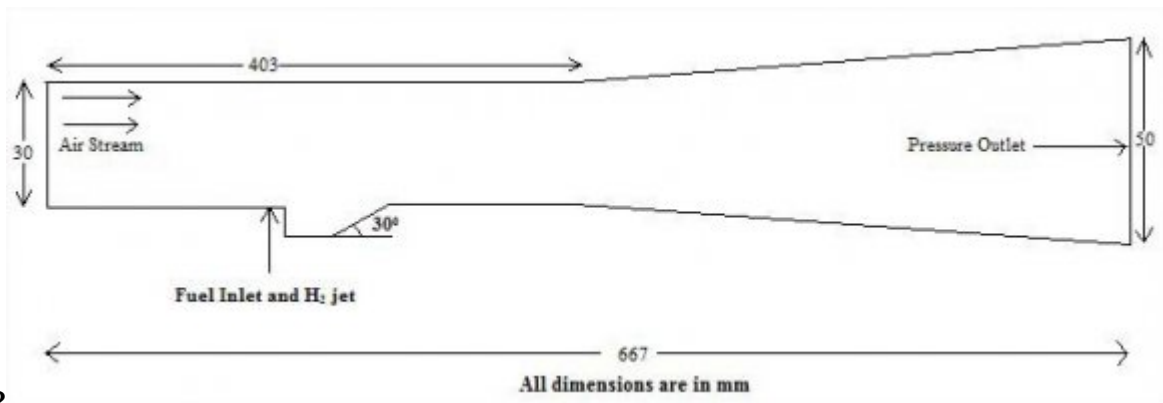
827

Figure 5: Fig 8 : 2 Fig 7 :



10

Figure 6: Fig 10 :



12

Figure 7: Fig 12 :

186 [Combustors ()] , Combustors . *Journal of Propulsion and Power* 2001. 17 (1) p. .

187 [Aiaa/Asme/Sae/ (2005)] , Aiaa/Asme/Sae/ . *ASEE Joint Prop. Conf* Jul 2005. p. .

188 [Yakar and Hanson (2001)] ‘Cavity Flame holders for Ignition and Flame Stabilization in Scramjets: An
189 Overview’. Ben Yakar , A Hanson , R . *Journal of Propulsion and Power* July-August 2001. 17 (4) p.
190 .

191 [Ben-Yakar and Hanson (1998)] ‘Cavity Flame holders for Ignition and Flame Stabilization in Scramjets: Review
192 and Experimental Study’. A Ben-Yakar , R Hanson . *AIAA Paper* July 1998. p. .

193 [Hsu et al. ()] ‘Characteristics of a Trapped-Vortex Combustor’. K.-Y Hsu , L P Goss , W M Roquemore . *Journal
194 of Propulsion and Power* 1998. 14 (1) p. .

195 [Jiyuntu et al. ()] *Computational Fluid Dynamics*, Guan Jiyuntu , Hengyeoh , Chaoqunliu . 2008. Elsevier Inc.

196 [Yu et al. (1998)] *Experimental Investigation on Dual-Purpose Cavity in Supersonic Reacting Flows*, K Yu , K J
197 Wilson , R A Smith , K C Schadow . Jan. 1998. (AIAA Paper 98-0723)

198 [Hsu et al. (2000)] *Fuel Distribution about a Cavity Flame holder in Supersonic Flow*, K.-Y Hsu , C Carter , J
199 Crafton , M Gruber , J Donbar , T Mathur , D Schommer , W Terry . July 2000. (AIAA Paper 2000-3583)

200 [Allen et al.] *Fuel-Air Injection Effects on Combustion in Cavity-Based Flame holders in a Supersonic Flow*, W
201 H Allen , P I King , Gruber . AIAA-2005-4105. p. 41.

202 [Gruber et al.] *Fundamental Studies of Cavity-Based Flameholder Concepts for Supersonic*, M R Gruber , R A
203 Baurle , T Mathur , K.-Y Hsu .

204 [Tishk Off et al. (1997)] ‘Future Direction of Supersonic Combustion Research: Air Force/NASA Workshop on
205 Supersonic Combustion’. J M Tishk Off , J P Drummond , T Edwards , A S Nejad . *AIAA Paper* Jan. 1997.
206 p. .

207 [Li et al. (2010)] Weipeng Li , Akira Takunomura , Kozo Oyama , Fujii . *LES Study of Feedback-loop
208 Mechanism of Supersonic Open Cavity Flows”; 40th Fluid Dynamics Conference and Exhibit*, 2010-5112,
209 28 June -1 July 2010. AIAA.

210 [Davis (1996)] ‘Numerical Analysis of Two and Three Dimensional Recessed Flame Holders for Scramjet
211 Applications’. D L Davis . *Ph.D. Dissertation, Aeronautics and Astronautics Dept., Air Force Inst. of
212 Technology* Sept. 1996.

213 [Baurle et al. (1998)] *Numerical and Experimental Investigation of a Scramjet Combustor for Hypersonic Missile
214 Applications*, R A Baurle , T Mathur , M R Gruber , Jackson , KR . July 1998. (AIAA Paper 98-3121)

215 [Pandey and Singh (2012)] ‘Numerical simulation of combustion chamber without cavity at Mach 3.12’. K M
216 Pandey , A P Singh . *International Journal of Soft Computing and Engineering (IJSCE)* 2231-2307. March
217 2012. (2) .

218 [Pandey and Reddy (2012)] ‘Numerical Simulation of Wall Injection with Cavity in Supersonic Flows of Scramjet
219 Combustion’. K M Pandey , S K Reddy , KK . *International Journal of Soft Computing and Engineering
220 (IJSCE)* 2231-2307. March 2012. (2) .

221 [Huang et al. (2010)] *Numerical Simulations of a Typical Hydrogen Fuelled Scramjet Combustor with a Cavity
222 Flame holder*, Wei Huang , Mohamed Shi-Bin Luo , Lin Pourkashanian , Derek B Ma , Jun Ingham , Zhen-Guo
223 Liu , Wang . WCE 2010. July 2010. London, UK.

224 [Eklund et al. (2001)] ‘Numerical Study of a Scramjet Combustor Fuelled by an Aerodynamic Ramp Injector in
225 Dual-Mode Combustion’. D R Eklund , R A Baurle , M R Gruber . *AIAA Paper* 2001-0379. Jan. 2001.

226 [Eklund and Gruber (1999)] ‘Study of a Supersonic Combustor Employing an Aerodynamic Ramp Pilot Injector’.
227 D R Eklund , M R Gruber . *AIAA Paper* June 1999. p. .

228 [Baurle and Gruber (1998)] *Study of Recessed Cavity Flow fields for Supersonic Combustion Applications*, R A
229 Baurle , M R Gruber . Jan. 1998. (AIAA Paper 98-0938)

230 [Katta and Roque More ()] ‘Study on Trapped-Vortex Combustor-Effect of Injection on Flow Dynamics’. V R
231 Katta , W M Roque More . *Journal of Propulsion and Power* 1998. 14 (3) p. .

232 [Mathur et al. (2001)] ‘Supersonic Combustion Experiments with a Cavity-based fuel injector’. T Mathur , M R
233 Gruber , Jackson K Donbar , J Donaldson , W Jackson , T Billig , F . *Journal of Propulsion and Power*
234 N0v-Dec 2001. 17 (6) p. .