

# On Dynamics of a Landing Gear Mechanism With Torsional Freeplay

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## Abstract

In this study, dynamics of a landing gear mechanism with torsional degree of freedom and torsional freeplay is analyzed. Derivation of the equations of motion of the model with torsional degree of freedom and the von Schlippe tire model are presented. Freeplay is introduced into the model and effects of freeplay angles of  $0^\circ$ ,  $0.5^\circ$ ,  $1^\circ$  and  $1.5^\circ$  are observed by obtaining time histories of the torsion angle and lateral tire deformation and limit cycles of the torsion angle. Amplitudes and frequencies of oscillations of the time histories of the torsion angle and lateral tire deformation are presented.

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*Index terms*— dynamics, torsional, Derivation, angles, histories, oscillations.

## 1 INTRODUCTION

ibration of aircraft steering systems has been a problem of great concern since the production of first airplanes. Shimmy is an oscillatory motion of the landing gear in lateral and torsional directions, caused by the interaction between the dynamics of the tire and the landing gear, with a frequency range of 10-30 Hz. Though it can occur in both nose and main landing gear, the first one is more common. Shimmy is a dangerous condition of selfexcited oscillations driven by the interaction between the tires and the ground that can occur in any wheeled vehicle. Problem of shimmy occurs in ground vehicle dynamics and aircraft during taxiing and landing. In other words, shimmy takes places either during landing, take-off or taxi and is driven by the kinetic energy of the forward motion of the aircraft. It is a combined motion of the wheel in lateral, torsional and longitudinal directions.

## 2 II.

## 3 SHIMMY

Shimmy can occur in steerable wheels of cars, trucks and motorcycles, as well as trailers and tea carts. In vehicle dynamics, shimmy is the unwanted oscillation of a rolling wheel about a vertical axis. It can occur in taxiing aircraft, as well. In the case of a shopping cart wheel, it is caused by the coupling between transverse and pivot degrees of freedom of the wheel. In the case of landing gear, shimmy is the result of the coupling between tire forces and landing gear bending and torsion. In other words, basic cause of shimmy is energy transfer from tire-ground contact force and vibration modes of the landing gear system.

Shimmy is an unstable phenomenon and it occurring with a certain combination of parameters such as mass, elastic quantities, damping, geometrical quantities, speed, excitation forces and nonlinearities such as friction and freeplay. It is difficult to determine shimmy analytically since it is a very complex phenomenon, due to factors such as wear and ground conditions that are hard to model. Small differences in physical conditions can lead to extremely different results. For example, it is reported in [1] that a new small fighter aircraft whose name is withheld, has displayed to vibrations during low and high speed taxi tests and first several landings and takeoffs, but shimmy vibrations with frequencies in the range 22-26 Hz were experienced during next several landings and take-offs at certain speeds, especially during landing. This demonstrates the effect of wear on landing gear shimmy. In the reported case, it was seen that tightening the rack too tight against the pinion prevented the

43 wheel from turning, while tightening it less tight caused the vibration to disappear but reappear in the following  
44 flights.

45 Ground control of aircraft is extremely important since severe shimmy can result in loss of control or fatigue  
46 failure of landing gear components. Vibration of aircraft steering systems deserves and has gained attention since  
47 shimmy is one of the most important problems in landing gear design. Shimmy is reported to be due to the forces  
48 produced by runway surface irregularities and nonuniformities of the wheels [2][3][4][5]. Modeling of aircraft tires  
49 presents similar challenges to those involved in modeling automotive tires in ground vehicle dynamics, on a much  
50 larger scale in terms size and loads on the tire [6].

51 Shimmy is a complex phenomenon influenced by many parameters. Causes of shimmy can be listed as follows  
52 [2],[7][8][9][10].

53 Insufficient overall torsional stiffness of the gear about the swivel axis Inadequate trail, since positive trail  
54 reduces shimmy Improper wheel mass balancing about the swivel axis Excessive torsional freeplay Low torsional  
55 stiffness of the strut Flexibilities in the design of the suspension Surface irregularities Nonuniformities of the  
56 wheels Worn parts

## 57 4 III. DETECTION AND SUPPRESSION OF SHIMMY

58 Shimmy is a great concern in aircraft landing gear design and maintenance. Prediction of nose landing gear  
59 shimmy is an essential step in landing gear design because shimmy oscillations are often detected during the taxi  
60 or runway tests of an aircraft, when it is no longer feasible to make changes on the geometry or stiffness of the  
61 landing gear. Although shimmy was observed in earlier aircraft as well, there were no extra shimmy damping  
62 equipments installed. Historically, France and Germany tended to deal with shimmy in the design phase, while  
63 in United States, the trend was to solve the problem after its occurrence. Currently, the general methodology  
64 is to employ a shimmy damper and structural damping. A shimmy damper, acting like a shock absorber in a  
65 rotary manner, is often installed in the steering degree of freedom to damp shimmy. It is a hydraulic damper with  
66 stroke limited to a few degrees of yaw. A shimmy damper restrains the movement of the nose wheel, allowing the  
67 wheel to be steered by moving it slowly, but not allowing it to move back and forth rapidly. It consists of a tube  
68 filled with hydraulic fluid causing velocity dependent viscous damping forces to form when a shaft and piston  
69 are moved through the fluid. Oleo-pneumatic shock absorbers are the most common shock absorber system in  
70 medium to large aircraft, since they provide the best shock absorption ability and effective damping. Such an  
71 absorber has two components: a chamber filled with compressed gas, acting as a spring and absorbing the vertical  
72 shock and hydraulic fluid forced through a small orifice, forming friction, slowing the oil and causing damping.  
73 Another common cure is to replace the tires even though they may not be worn out [10][11][12].

74 Shimmy started being investigated in 1920's both theoretically and experimentally and soon it became clear  
75 that it is caused not by a single parameter but by the relationships between parameters. Effects of acceleration  
76 and deceleration on shimmy have been reported to be examined, and the accelerating system is found to be  
77 slightly less stable [13]. Number of publications available in literature on landing gear shimmy is limited because  
78 many developments are proprietary and are not published in literature.

79 IV.

## 80 5 LITERATURE SURVEY

81 Many papers have been published addressing shimmy as a vehicle dynamics problem. In that perspective, tire is  
82 the most important item, and tire models have been investigated. [13] examines the wheel shimmy problem and  
83 its relationship with longitudinal tire forces, vehicle motions and normal load oscillations. [8] compares different  
84 dynamic tire models for the analysis of shimmy instability. [3] is an investigation of tire parameter variations  
85 in wheel shimmy, by considering the shimmy resulting from the elasticity of a pneumatic tire, particularly in  
86 taxiing aircraft. [14] is on the application of perturbation methods to investigate the limit cycle amplitude and  
87 stability of the wheel shimmy problem. [7] deals with the shimmy stability of twin-wheeled cantilevered aircraft  
88 main landing gear. The objective in [15] is to develop software on assessing shimmy stability of a general class of  
89 landing gear designs using linear and nonlinear landing gear shimmy models. [16] studies the periodic shimmy  
90 vibrations and chaotic vibrations of a simplified wheel model using bifurcation theory. [17] is on tire dynamics  
91 and is a development to deal with large camber angles and inflation pressure changes. [18] is another study on  
92 tire dynamics, where stability charts show the behavior of the system in terms of certain parameters such as  
93 speed, caster length, damping coefficient and relaxation length. [19] is an experimental study on wheel shimmy  
94 where system parameters are identified, stability boundaries and vibration frequencies are obtained on a test rig  
95 for an elastic tire. Dependence of shimmy oscillations in the nose landing gear of an aircraft on tire inflation  
96 pressure are investigated in [20]. The model derived in [21] is used and it is concluded that landing gear is less  
97 susceptible to shimmy oscillations at inflation pressures higher than the nominal.

98 Transverse vibrations of landing gear struts with respect to a hull of infinite mass have been studied  
99 theoretically in [22]. Similarly, [23] presents a nonlinear model describing the dynamics of the main gear wheels  
100 relative to the fuselage.

101 Lateral dynamics of nose landing gear shimmy models has gained some attention. Lateral response of a nose  
102 landing gear has been investigated in [10] where nonlinearities arise due to torsional freeplay. In [24], lateral

103 response to ground-induced excitations due to runway roughness is taken into consideration as well. Lateral  
 104 stability of a nose landing gear with a closed loop hydraulic shimmy damper is presented in [12]. Closed form  
 105 analytical expressions for shimmy velocity and shimmy frequency are derived in regard to the lateral dynamics  
 106 of a nose landing gear in [25].

107 A dynamic model of an aircraft nosegear is developed in [9] and effects of design parameters such as energy  
 108 absorption coefficient of the shimmy damper, the location of the center of gravity of the landing gear, shock strut  
 109 elasticity, tire compliance, friction between the tire and the runway surface and the forward speed on shimmy  
 110 are investigated. It is shown in [26] that dry friction is one of the principal causes of shimmy.

111 Bifurcation analysis of a nosegear with torsional and lateral degrees of freedom is performed in [21]. Similarly,  
 112 bifurcation analysis of a nosegear with torsional, lateral and longitudinal modes is performed in [27]. In a more  
 113 mathematical study, incremental harmonic balance method is applied to an aircraft wheel Theoretical research  
 114 on shimmy has a long history, with the initial focus on tire dynamic behavior because tires play an important  
 115 role in causing shimmy instability. Theories on tire models can be divided into stretched string models and point  
 116 contact models. In the stretched string model proposed by von Schlippe, the tire centerline is represented as a  
 117 string in tension, the tire sidewalls are represented by a distributed spring where the string rests and the wheel  
 118 is represented by a rigid foundation for the spring. Pacejka has proposed replacing the string by a beam. The  
 119 point contact method assumes the effects of the ground on the tire act at a single contact point and is much  
 120 easier to implement in an analytical model.

121 V. where  $I_z$  is the moment of inertia about the  $z$  axis,  $M_1$  is the linear spring moment between the turning  
 122 tube and the torque link,  $M_2$  is the combined damping moment from viscous friction in the bearings of the  
 123 oleo-pneumatic shock absorber and from the shimmy damper, (2)

## 124 6 MATHEMATICAL MODE

125 (3) (4) (5) where  $k$  is the torsional spring rate,  $c$  is the torsional damping constant,  $v$  is the taxiing velocity and  
 126  $k$  is the tread width moment constant defined as [29]  $k M = 1 ? c M = 2 y z e F M M ? = 3 ? ? v M = 4 z F$   
 127  $F c a ? ? 2 15 . 0 ? = (6)$

128  $F_y$  and  $M_z$  depend on the vertical force  $F_z$  and slip angle  $\delta$ . Tire sideslip characteristics are nonlinear.  
 129 Cornering force  $F_y$  and vertical force  $F_z$  are related as (7) (8)

130 Where  $\delta_{lim}$  is the limiting slip angle or the limit angle of tire force and sign is the sign function defined as (9) Slip  
 131 angle may be caused by either pure yaw or pure sideslip. Pure yaw occurs when the yaw angle is allowed to vary  
 132 while the lateral deflection  $y$  is held at zero. Pure sideslip, on the other hand, occurs when the lateral deflection  
 133  $y$  is allowed to vary as the yaw angle is held at zero [11].

134 An expression is given for the nonlinear sideslip characteristic in the widely used Magic Formula [7,11,17] as  
 135 the following (10) where  $B, C, D$  and  $E$  are functions of the wheel load, slip angle, slip ratio and camber.  $B$  and  
 136  $E$  are related to vertical force  $F_z$ ,  $C$  is the shape factor and  $D$  is the peak value of the curve.

137 Plots of  $F_y$  versus  $F_z$  will not be presented here due to lack of space, but they have similar characteristics when  
 138 obtained using either (7) and (8) or the Magic Formula, thus the simple approximations given by (??) and (??)  
 139 are used instead of the complicated Magic Formula. Only force and moment derivatives are needed as parameters  
 140 for (7) and (8).

141 Aligning moment  $M_z$  is defined using a halfperiod sine.

142  $M_z$  is approximated by a sinusoidal function and the constant zero given by (??1) and (??2). (??1) (12) where  
 143  $\delta_{lim}$  is the limiting angle of tire moment.

144 b) Tire model Tire is modeled using the elastic string theory. Lateral deflection of the tire is described as  
 145 [11, ??9] (13) Ground forces are transmitted to the wheel through the tire, and these forces acting on the tire  
 146 footprint deflect the tire. Elastic string theory states that lateral deflection  $y$  of the leading contact point of the  
 147 tire with respect to tire plane can be described as a first order differential equation given by (13). This equation  
 148 is derived as follows.

149 Tire sideslip velocity  $V_t$  is expressed as (14) Where  $\tau$  is the time constant,  $l$  is the relaxation length, which is  
 150 the ratio of the slip stiffness to longitudinal force stiffness. The tire also undergoes yaw motion, leading to a yaw  
 151 velocity  $V_r$  which is approximated as Global Journal of Researches in Engineering Volume XII Issue v v v v I  
 152 Version I 38 ( D D D D ) D 2012 ebruary F ? ? F z y c F F = , for ? ? ? ( ) ? ? ? sign c F F F z y = , for ? ?  
 153  $> ? ( ) ? ( ) ? ? > = ? ? ? ? ?$  if, 1 if, 1 sign ? ? ( ) ( ) { } [ ] ? ? ? B B E B C D F y arctan arctan sin ? ?  
 154  $= z y F F ? z z F M = ? ? ? ? g g M z z c F M 180 \sin 180$  , for  $g ? ? ? 0 = z z F M$  , for  $g ? ? > g ? ( ) ? ? ?$   
 155  $a e v y v y ? + = + ? y y V t + = V ? ? = ? ( ) ? ? a e v V r ? + = (15)$

156 As the wheel rolls on the ground, (??6) Substituting (??4) and (??5) into (??6) yields (??3).

157 An equivalent side slip angle caused by lateral deflection is used to compute cornering force  $F_y$  and aligning  
 158 moment  $M_z$  and is approximated as  $r_t V V = (17)$

159 Equations (??), (??3) and (??7) constitute the governing equations of the torsional motion of the landing  
 160 gear and include nonlinear tire force and moment. Parameters of a light aircraft used in the computations are  
 161 given in table ???.? ? ? y = ? arctan

162 Table ???: Parameters used in the torsional dynamics.



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219 Mathematical analysis of the behavior of a two dimensional aeroelastic system with freeplay nonlinearity is  
220 presented in [43] and two formulations are developed. Formulations are extended for a hysteresis model in [44].  
221 Unlike a freeplay model which consists of three linear subsystems, a hysteresis model consists of six.

222 Bifurcation analysis an airfoil having two degrees of freedom with both freeplay and cubic stiffness nonlinearities  
223 in pitch placed in supersonic flow has been conducted in [35]. Bifurcation analysis of an aircraft with freeplay  
224 nonlinearity is conducted in [45]. Limit cycle oscillations of an airfoil with two degrees of freedom having freeplay in  
225 the pitching degree of freedom are examined experimentally and theoretically in [34]. An experimental delta wing  
226 model with freeplay at the attachment points is designed and tested in [46], and its gust response is investigated  
227 in [47]. Effect of freeplay on the aerodynamic response, such as limit cycle flutter, has been examined. It has  
228 been found that the amplitude and position of the limit cycle varies with the magnitude of freeplay. Effects of  
229 variations in parameters have been examined for both the damped and limit cycle oscillations. Critical flutter  
230 speeds are predicted.

231 Hinges of control surfaces often demonstrate freeplay nonlinearity. [48] is a study examining the limit cycle  
232 oscillations of a combination of an airfoil and an aileron, resulting in three degrees of freedom, with freeplay in the  
233 aileron hinge. Aeroelastic response of other two dimensional systems having control surface freeplay nonlinearity  
234 are studied using the harmonic balance approach in [49] and both numerically and experimentally in [39]. A  
235 dissertation was presented to Duke University in 2000, covering the dynamics of a two dimensional aeroelastic  
236 system with control surface freeplay nonlinearity, both experimentally and mathematically [50]. Limit cycle  
237 oscillations are observed. The system is very similar to the one given in [48], a combination of an airfoil with an  
238 aileron.

239 A three dimensional control surface with play is investigated in [51] to demonstrate the effects of angle of  
240 attack and Mach number. Flutter analysis of a missile wing having freeplay in the rotation degree of freedom  
241 of the wing control mechanism is conducted in [33] by investigating limit cycles and chaotic motion. Results state  
242 that the system response depends on the amount of freeplay and initial conditions.

243 A study on a mechanical system exhibiting freeplay nonlinearity is studied both numerically and experimentally  
244 in [36] where the problem of developing a mathematical model and performing a simulation of the dynamics of  
245 systems exhibiting freeplay nonlinearity is addressed. Contact due to freeplay is considered, constraints are  
246 formulized and the stability of an aircraft wing displaying freeplay in the hinge supporting a control surface is  
247 investigated. Freeplay is considered as one of the rotor faults in the simulation of helicopter structural damages  
248 in [52].

249 Freeplay model used in this study is based on the ones in [31] and [38]. Dynamics of a landing gear mechanism  
250 with freeplay in the torsional degree of freedom is analyzed in [38], while dynamic behavior of a two dimensional  
251 airfoil with freeplay in pitch, oscillating in pitch and plunge directions, subjected to inviscid, transonic flow is  
252 analyzed in [31]. Both freeplay nonlinearities are modeled using the same principle and formulation, although  
253 the two studies are in two very distinct disciplines. Same formulation as in [31] is employed in [40, 53], and  
254 mathematical models given in [32,33,35,37,[41][42][43]48] are also similar .

255 Freeplay is modeled as a nonlinear spring as in figure 2, where some deflection is possible before a force develops  
256 and the spring force is zero if the amplitude remains within the freeplay band. Formulations have been suggested  
257 in literature to determine an equivalent linear stiffness. Equation 34 gives the piecewise continuous restoring  
258 moment function similar to the one used in [38] to describe the concentrated nonlinearity at the torsional degree  
259 of freedom.

$$260 \quad (29)$$

261 Torsion is denoted by  $\theta$  is the stiffness coefficient and  $\alpha$  is the freeplay angle.

## 262 12 VIII. INCORPORATION OF FREEPLAY INTO THE 263 LANDING GEAR MODEL

264 Torsional freeplay is incorporated into the equations of motion of the landing gear. Results are displayed for  
265 various values of the freeplay angle within the range  $0^\circ$ - $2^\circ$ , as this is the range employed in literature. Freeplay  
266 has been incorporated into the equations of motion of landing gear mechanisms in very few studies literature  
267 [38].

268 Freeplay model given in (29) can be incorporated in the equations of motion in two ways. One of them, is  
269 to linearize the model as in (53)-(28) and substitute (29) into in (23). This way, the only nonlinearity in the  
270 model is freeplay nonlinearity such that the second equation in (53) becomes (30) Second way of incorporating  
271 freeplay nonlinearity in the model is to obtain a more realistic model by substituting (29) directly into the  
272 nonlinear model. This is the approach taken here. Nonlinear equations are integrated using the fourth order  
273 Runge-Kutta algorithm.

## 274 13 IX. RESULTS

275 Effects of freeplay are observed by obtaining time histories of the torsion angle and lateral tire deformation and  
276 limit cycles. Freeplay angles of  $0^\circ$ ,  $0.5^\circ$ ,  $1^\circ$  and  $1.5^\circ$  are incorporated. Amplitudes and frequencies of oscillations  
277 of the time histories of the  $\theta$  ) Effect of freeplay on the torsion angle and lateral tire deformation are observed.  
278 By observing tables 2-4 it can be stated that the existence of a freeplay angle prevents shimmy damping of the

279 system with the same physical parameters. The increase in the freeplay angle increases shimmy amplitude. A  
280  $0.5^\circ$  increase of the freeplay angle from  $0.5^\circ$  to  $1^\circ$  doubles the amplitude in all 3 cases. Another  $0.5^\circ$  increase in  
281 the freeplay angle from  $1^\circ$  to  $1.5^\circ$  causes a 25% increase in the amplitude of the torsion angle and a 55% increase  
282 in the amplitude of the lateral tire deformation. ( ) ? ? + ? ? ? ? ? = fp fp fp fp fp fp K K M ? ? ? ? ? ? ? ?  
283 ? ? ? ? ? ? if if 0 if ? , ? K fp ? fp ? ? 1 c ( )

284 **14 Global Journal of Researches in Engineering**



Figure 1: F

285 1 2

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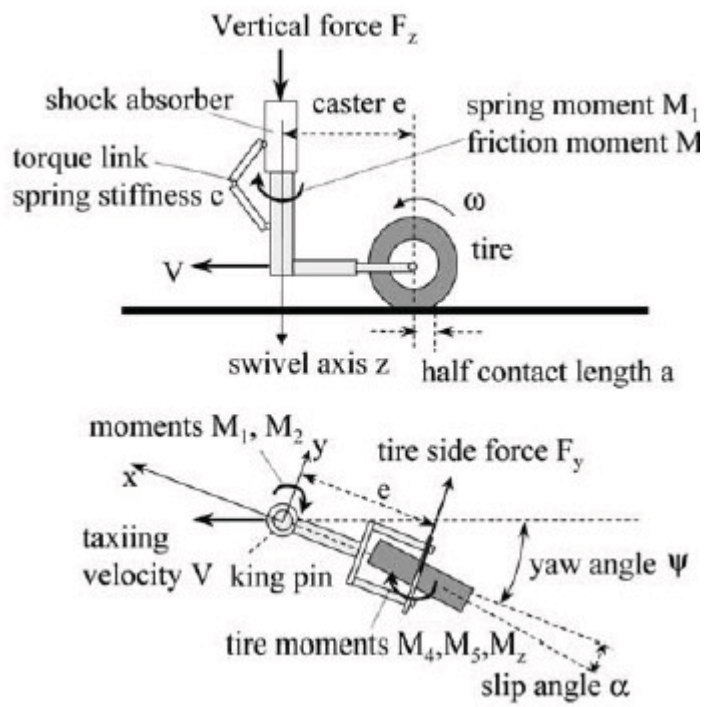


Figure 2:

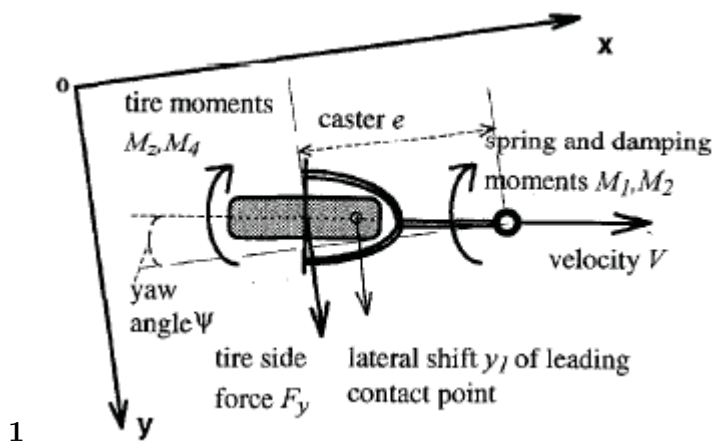


Figure 3: Figure 1 :

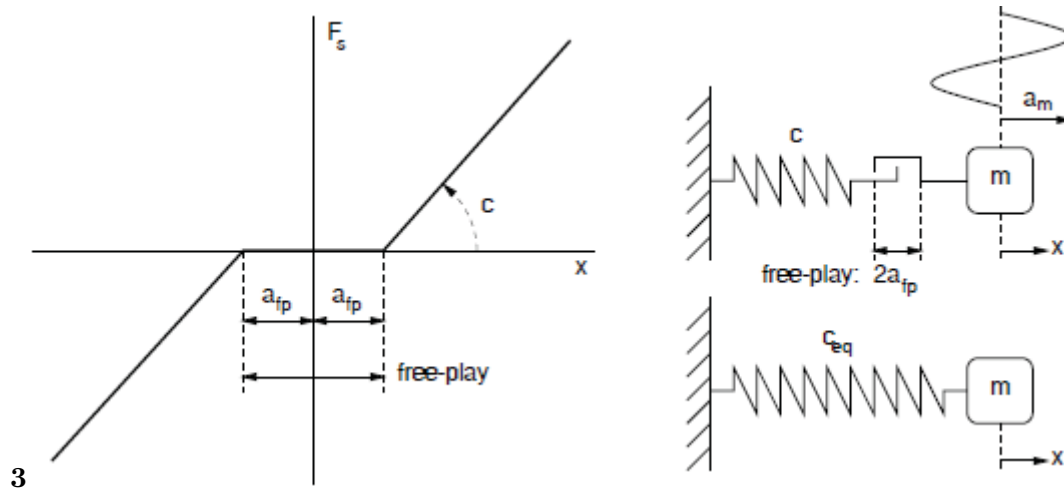
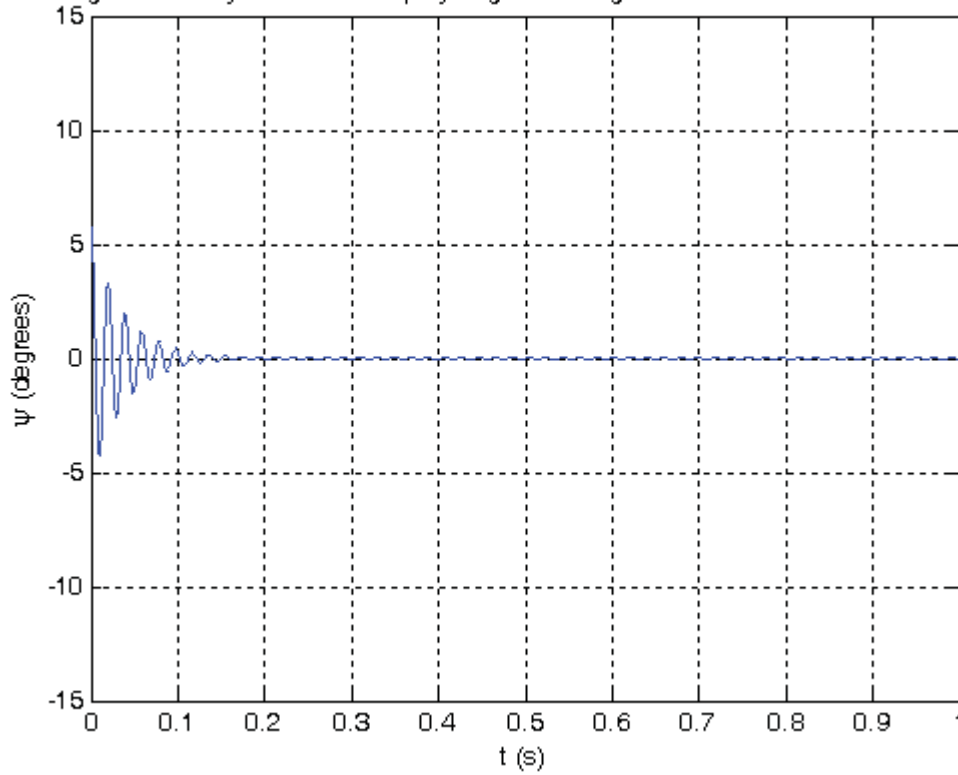


Figure 4: M 3

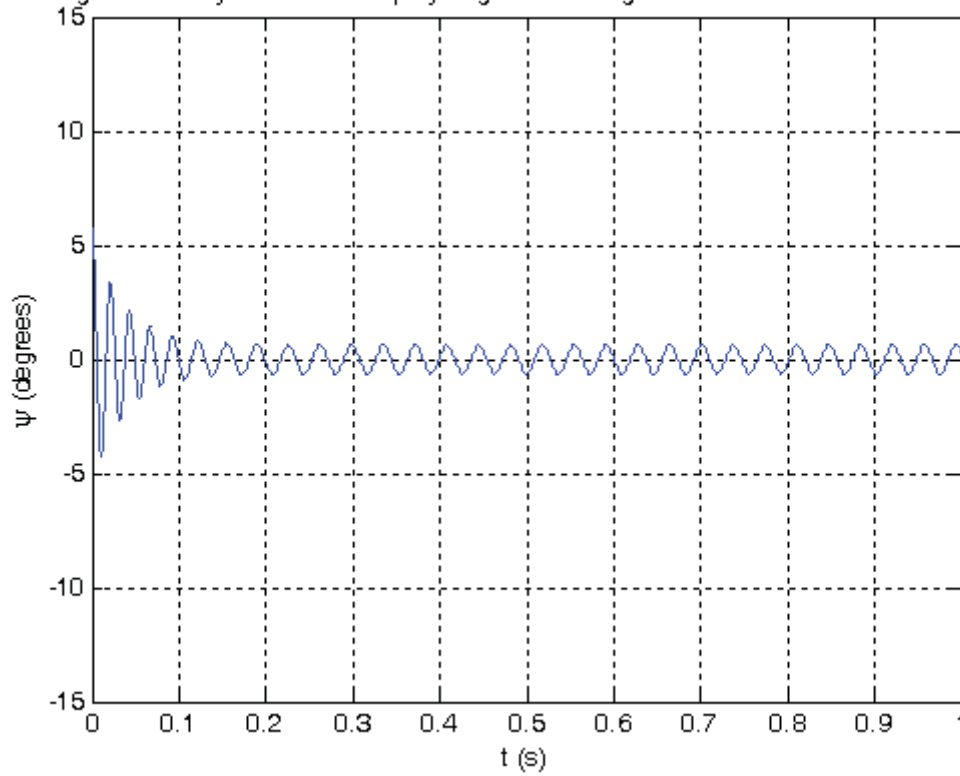
Torsion angle for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-100$  Nm/rad/



20

Figure 5: ( 20 )

Torsion angle for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-100$  Nm/rad



3

Figure 6: M 3

Torsion angle for the system with freeplay angle of 1 degree for  $v=50$  m/s and  $k=-100$  Nm/rad

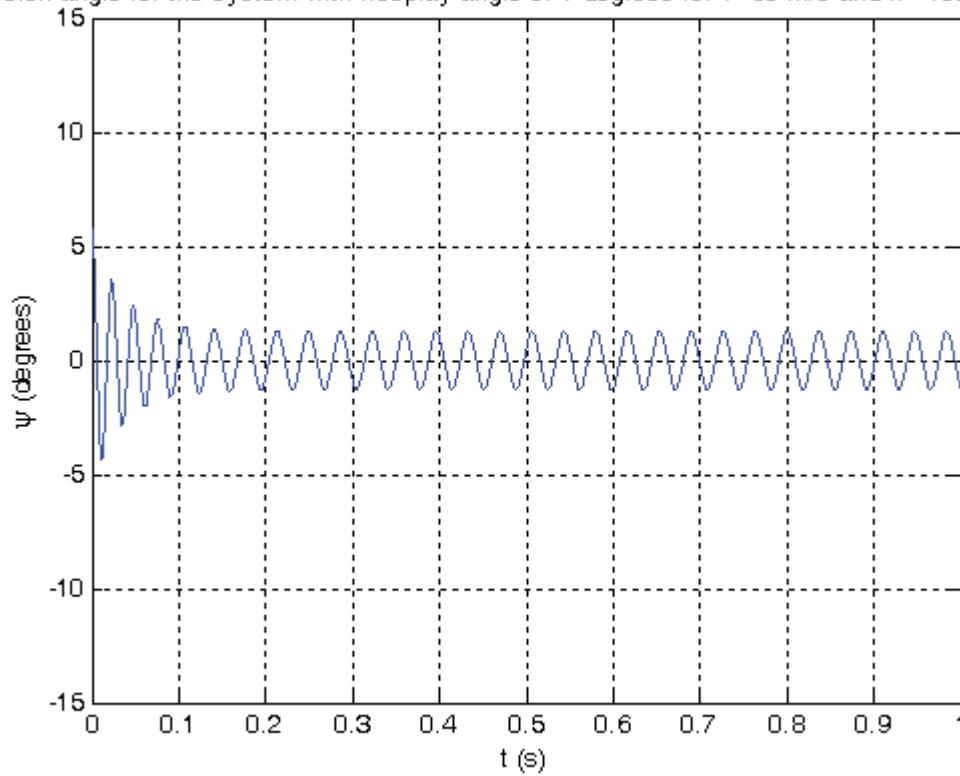
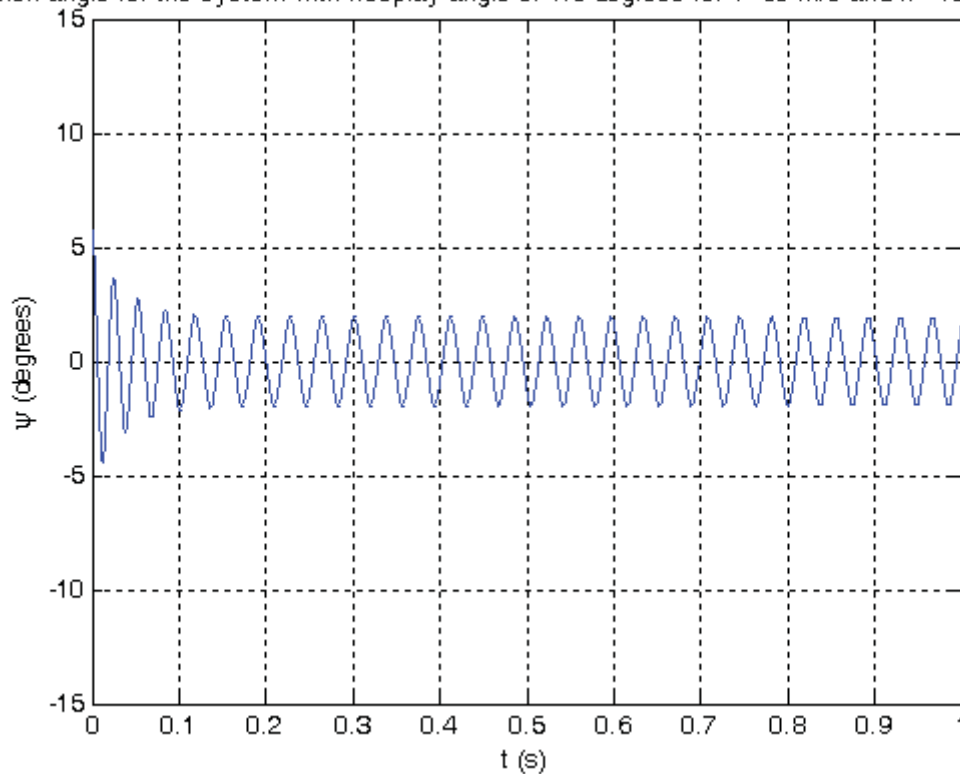


Figure 7: Volume

Torsion angle for the system with freeplay angle of 1.5 degrees for  $v=50$  m/s and  $k=-100$  Nm/rad



2

Figure 8: FFigure 2 :

lateral tire deformation for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-100$  N/r

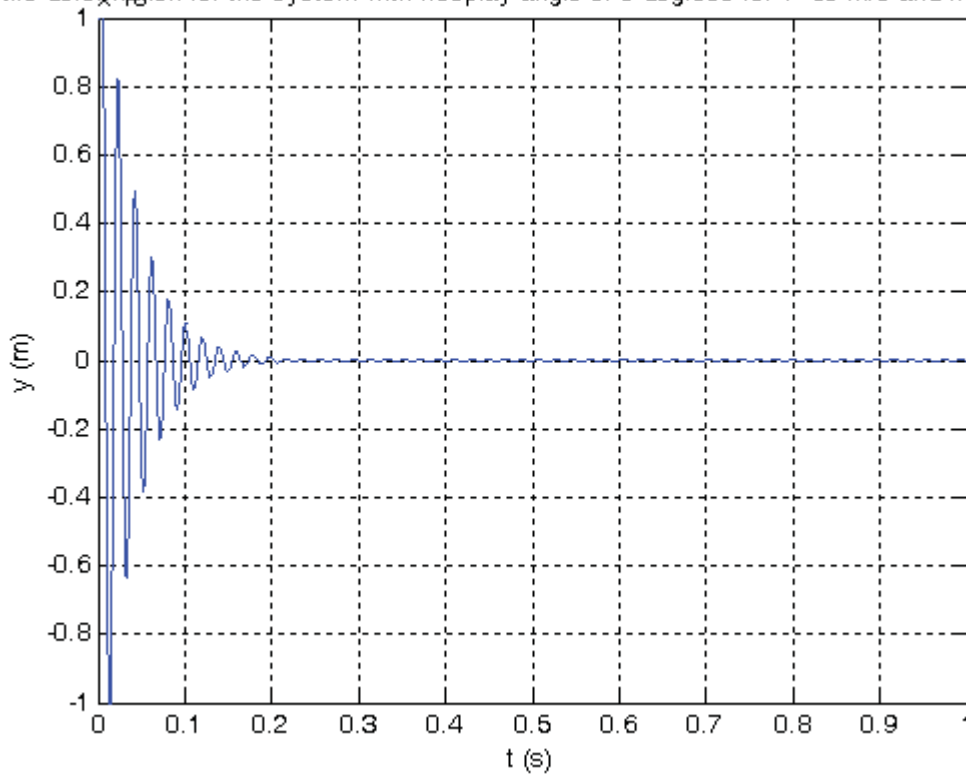


Figure 9:

lateral tire deformation for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-100$  N/r

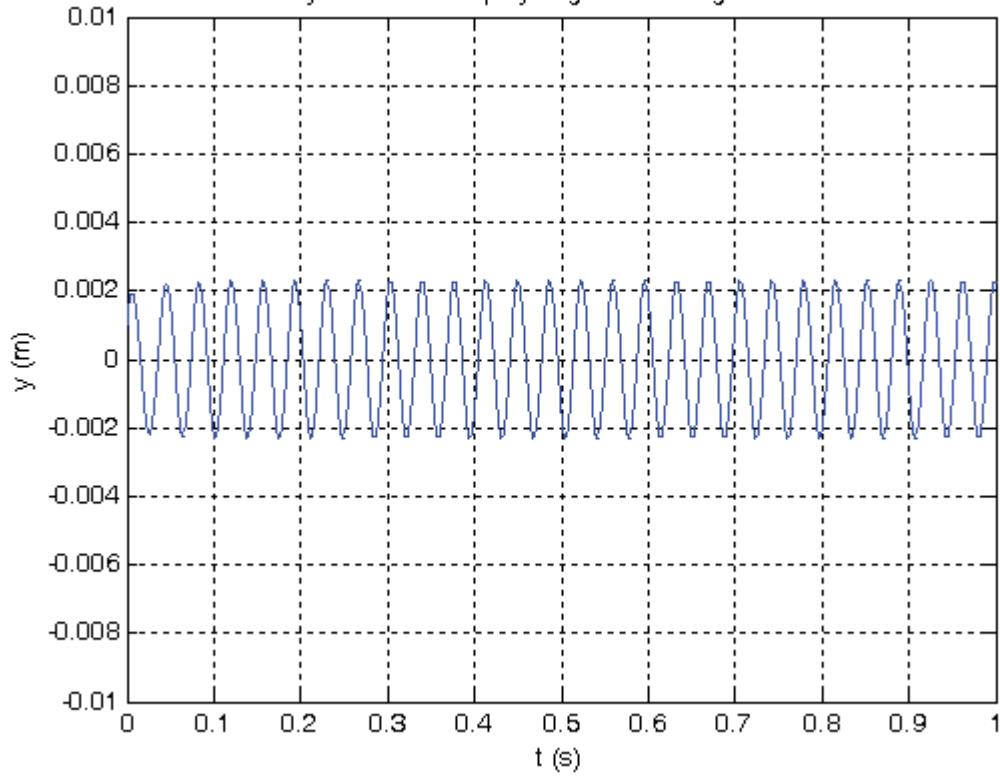
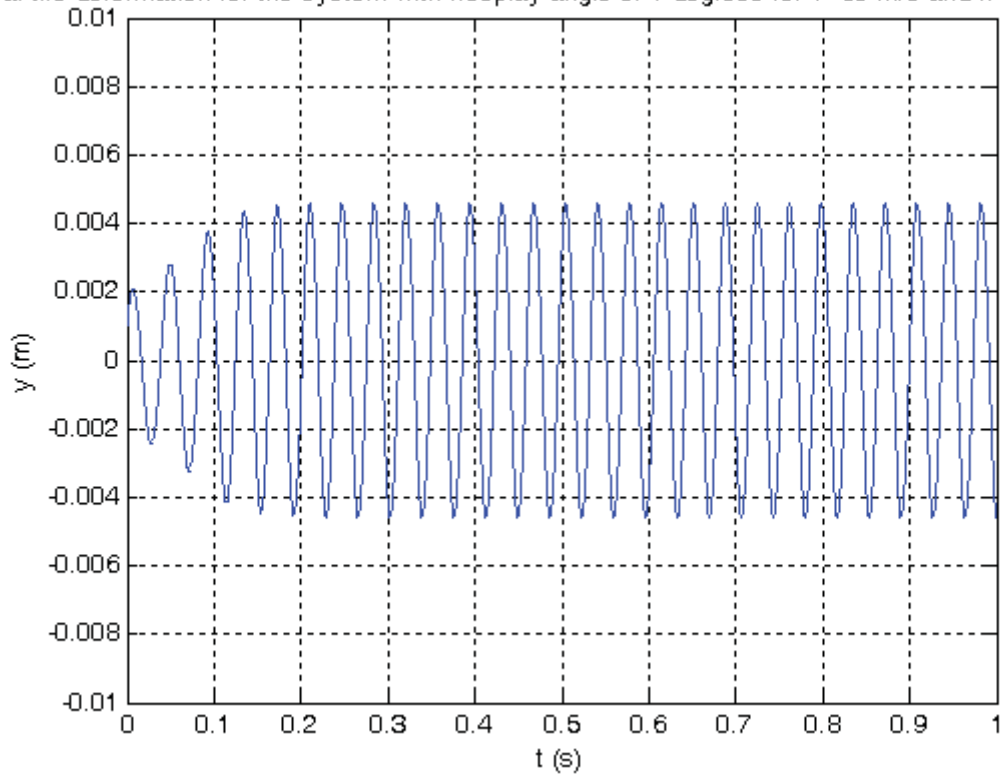


Figure 10:

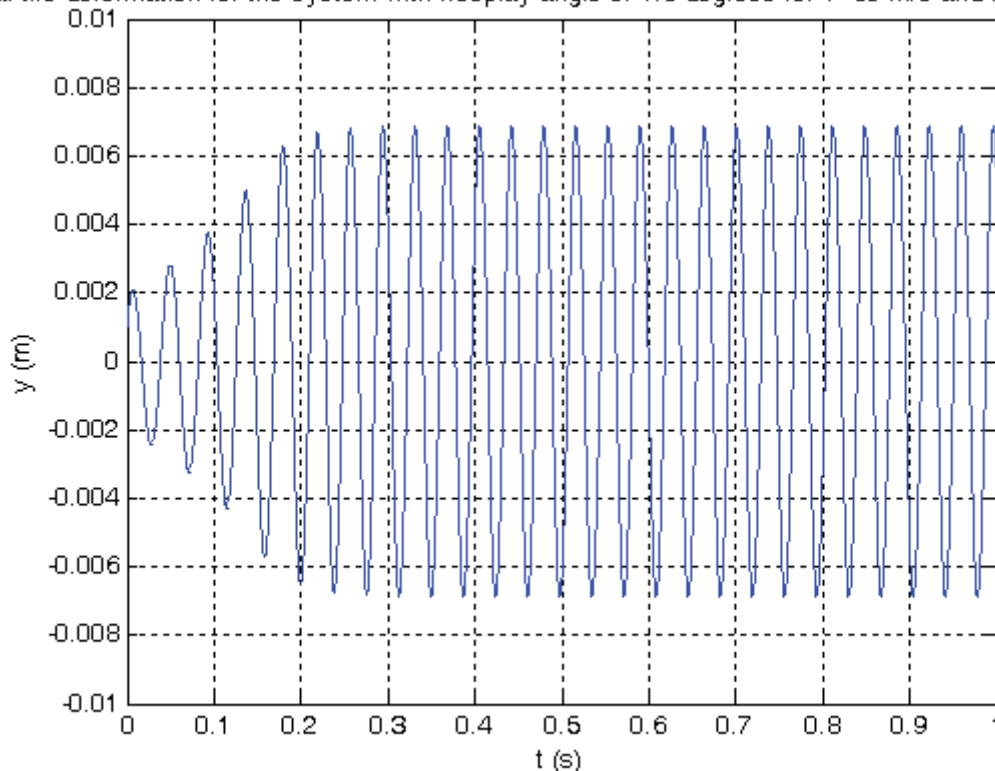
lateral tire deformation for the system with freeplay angle of 1 degrees for  $v=50$  m/s and  $k=-100$  N/r



3

Figure 11: Figure 3 :

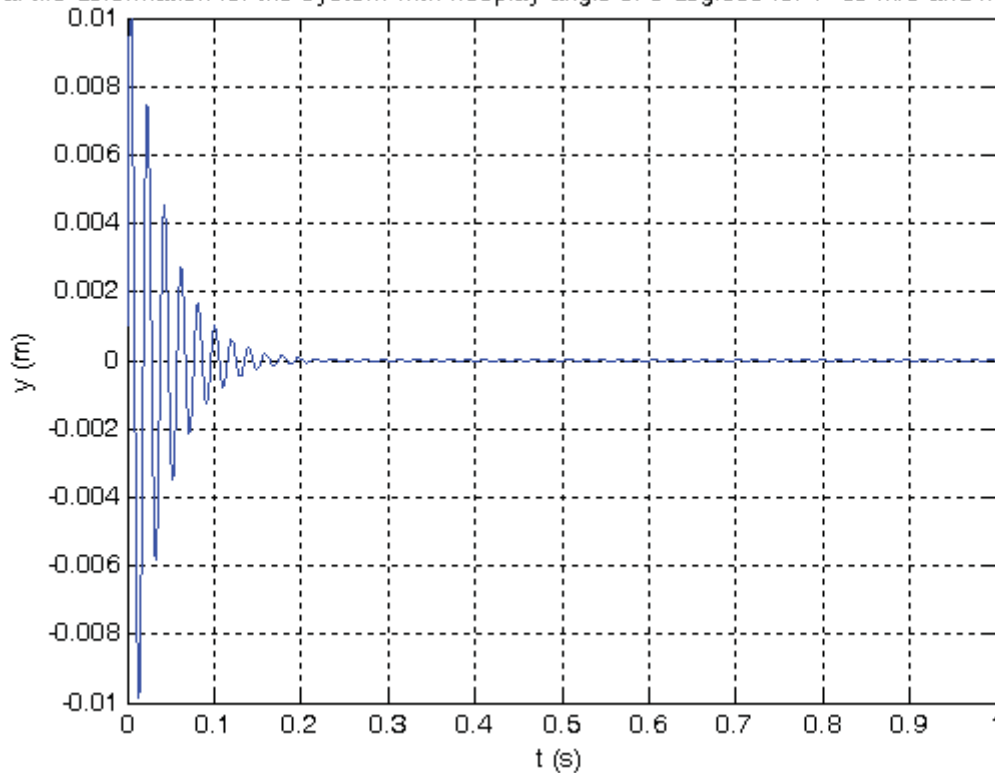
lateral tire deformation for the system with freeplay angle of 1.5 degrees for  $v=50$  m/s and  $k=-100$  N/r



4

Figure 12: Figure 4 :

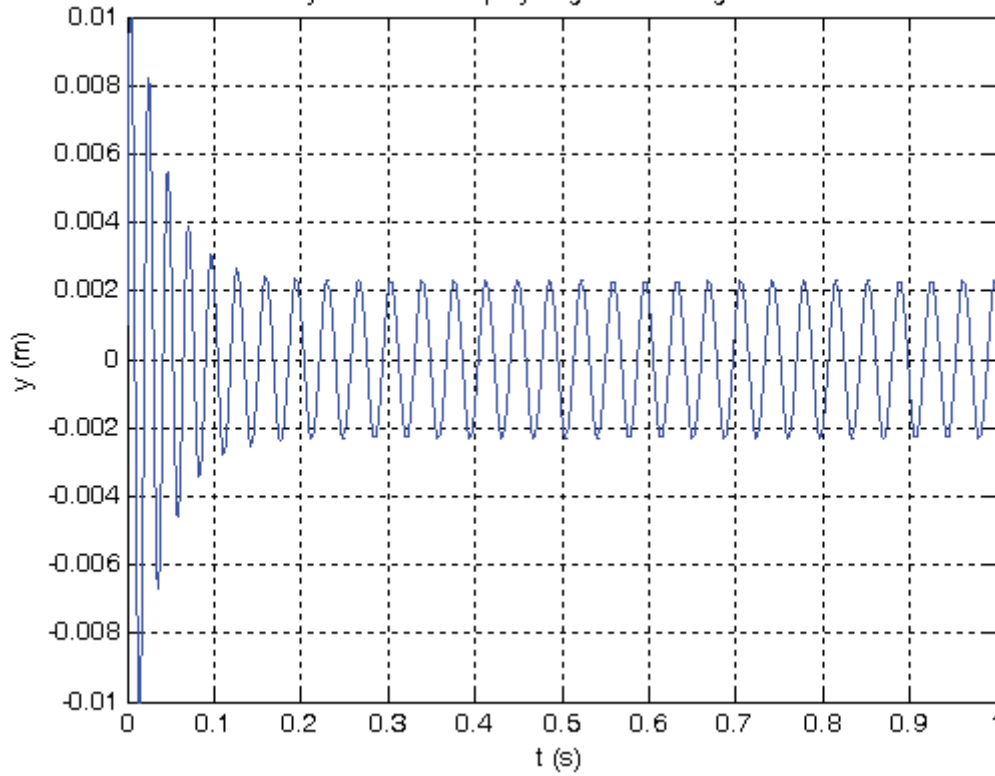
lateral tire deformation for the system with freeplay angle of 0 degrees for  $v=60$  m/s and  $k=-100$  N/r



5

Figure 13: Figure 5 :

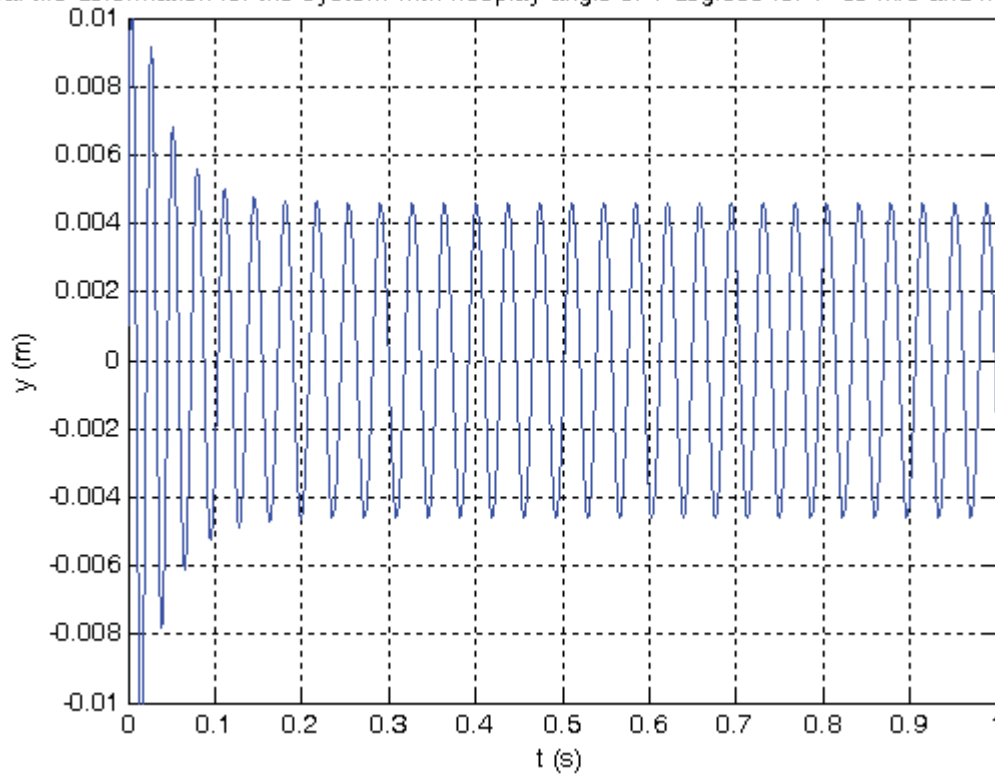
lateral tire deformation for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-100$  N/r



6

Figure 14: Figure 6 :

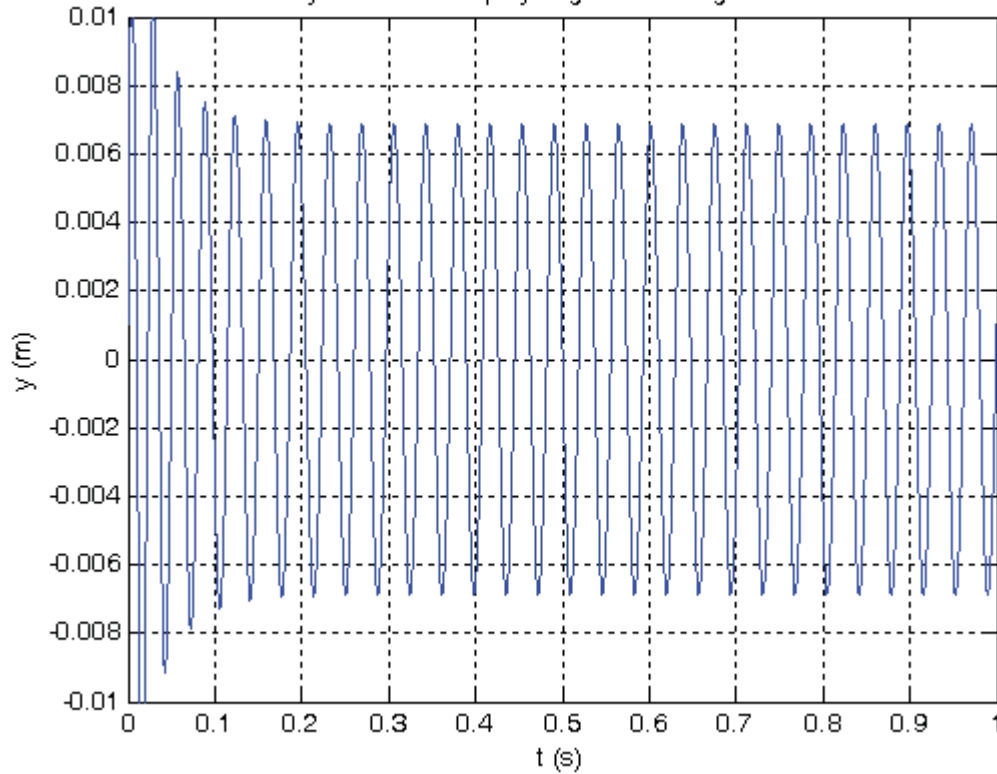
lateral tire deformation for the system with freeplay angle of 1 degrees for  $v=50$  m/s and  $k=-100$  N/r



27

Figure 15: Table 2 :Figure 7 :

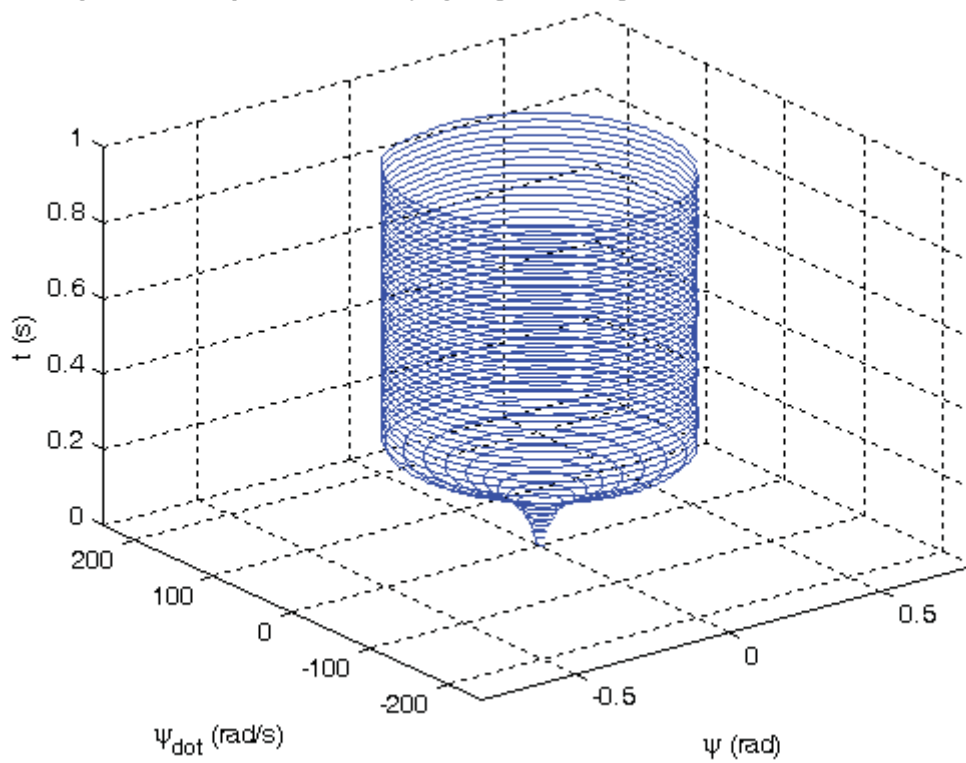
Vertical tire deformation for the system with freeplay angle of 1.5 degrees for  $v=50$  m/s and  $k=-100$  N/rad



89

Figure 16: Figure 8 :Figure 9 :

Limit cycle for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



3

Figure 17: Table 3 :

Limit cycle for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s

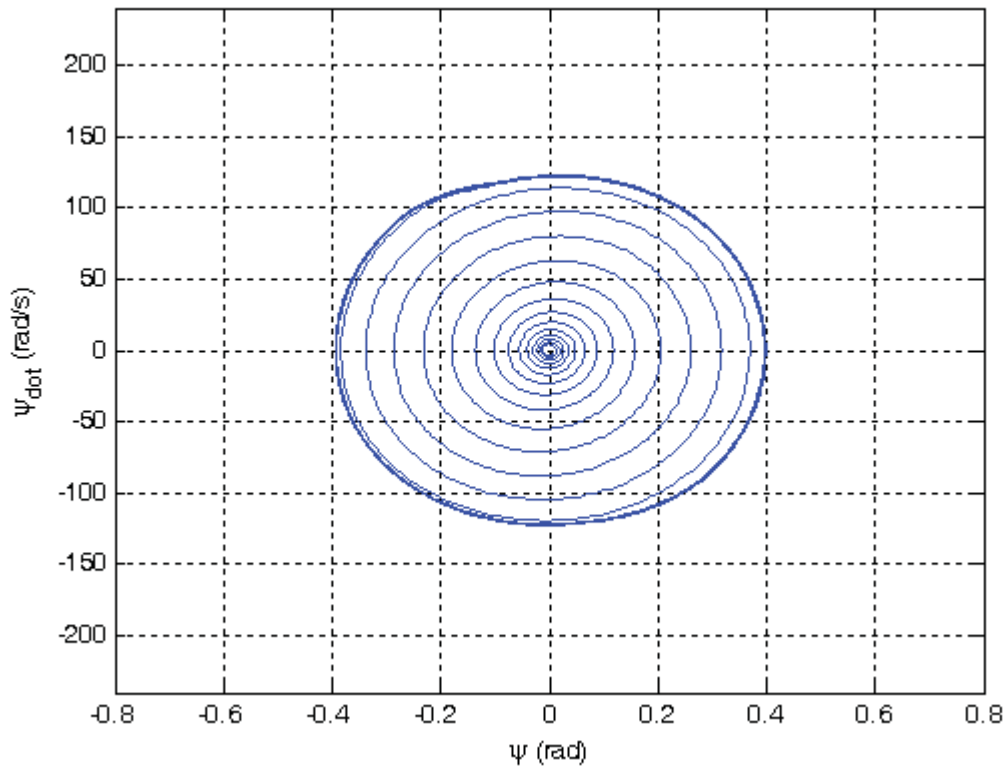
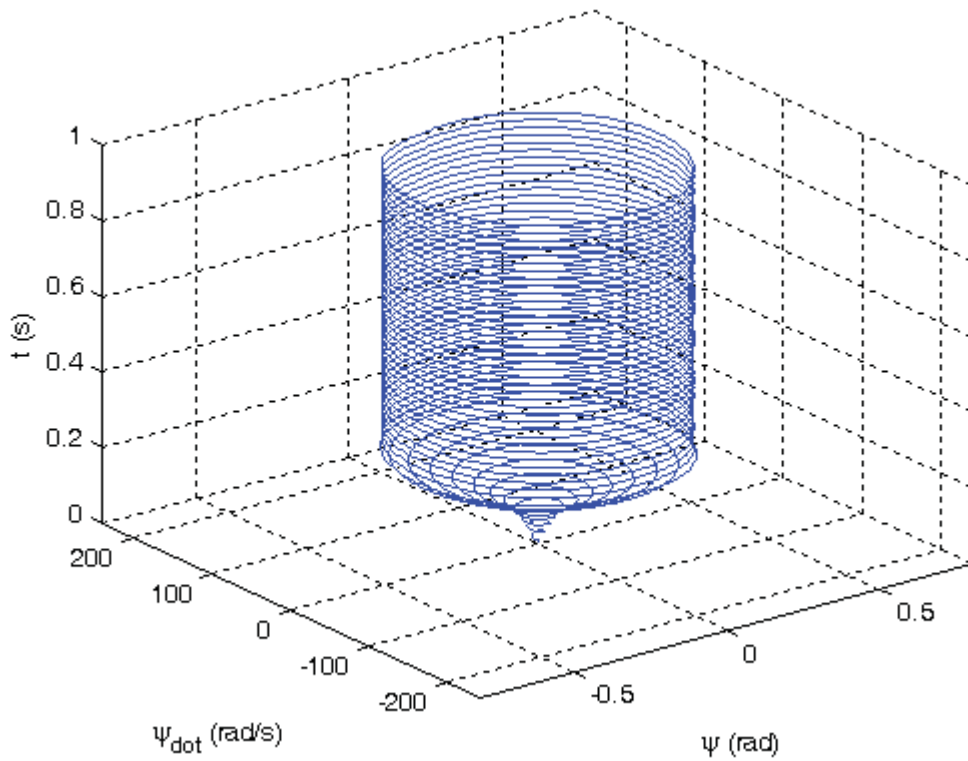


Figure 18: Figure11:

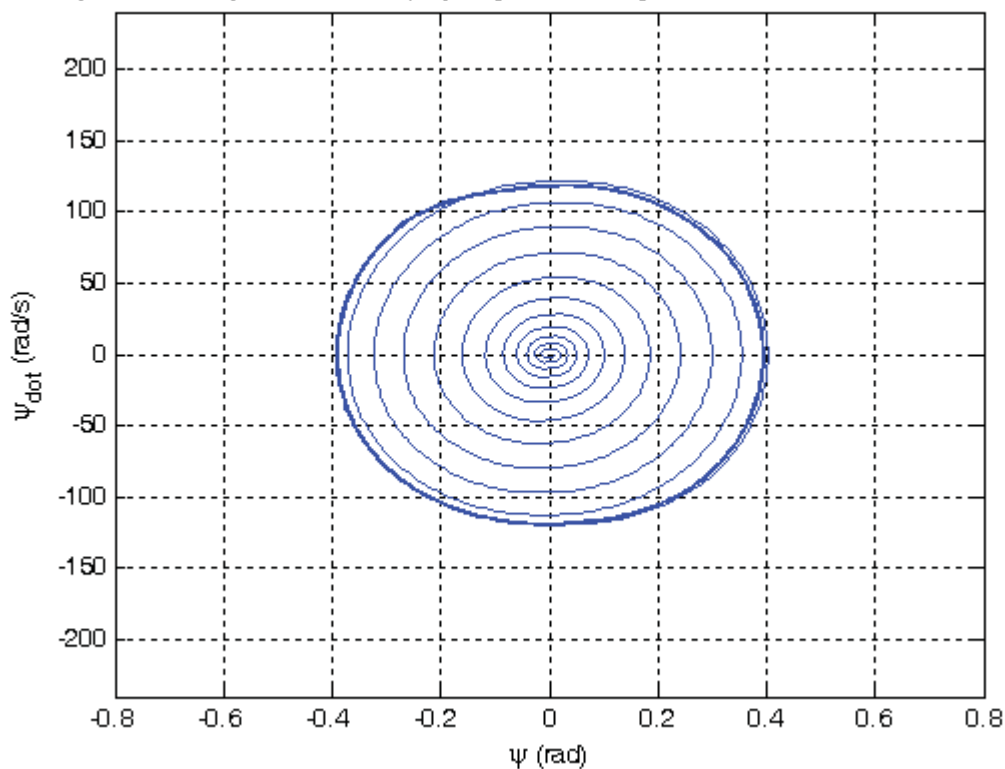
Limit cycle for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



124

Figure 19: Figure 12 :Figure13:Table 4 :

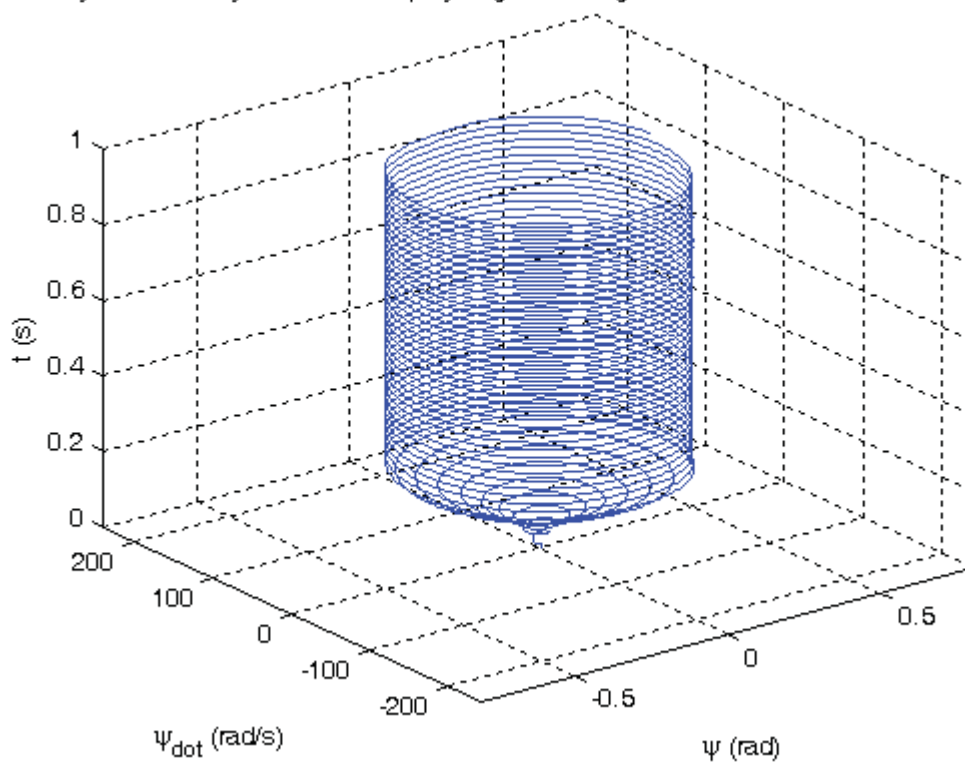
Limit cycle for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



15

Figure 20: Figure 15 :

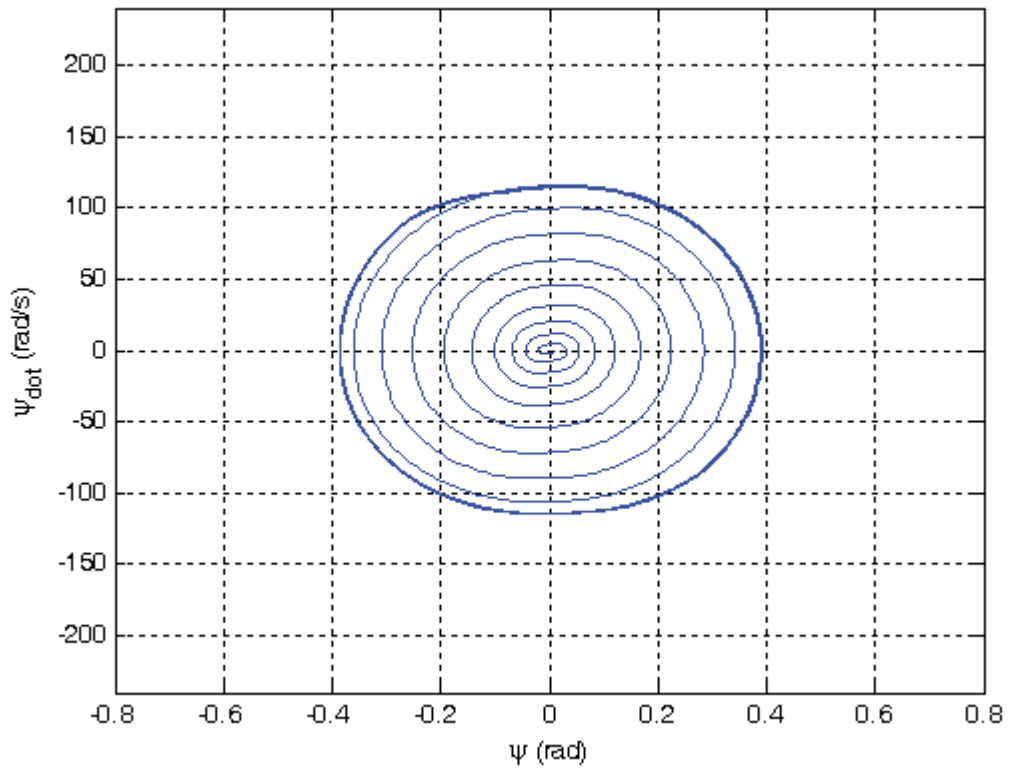
Limit cycle for the system with freeplay angle of 1 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



16

Figure 21: Figure 16 :

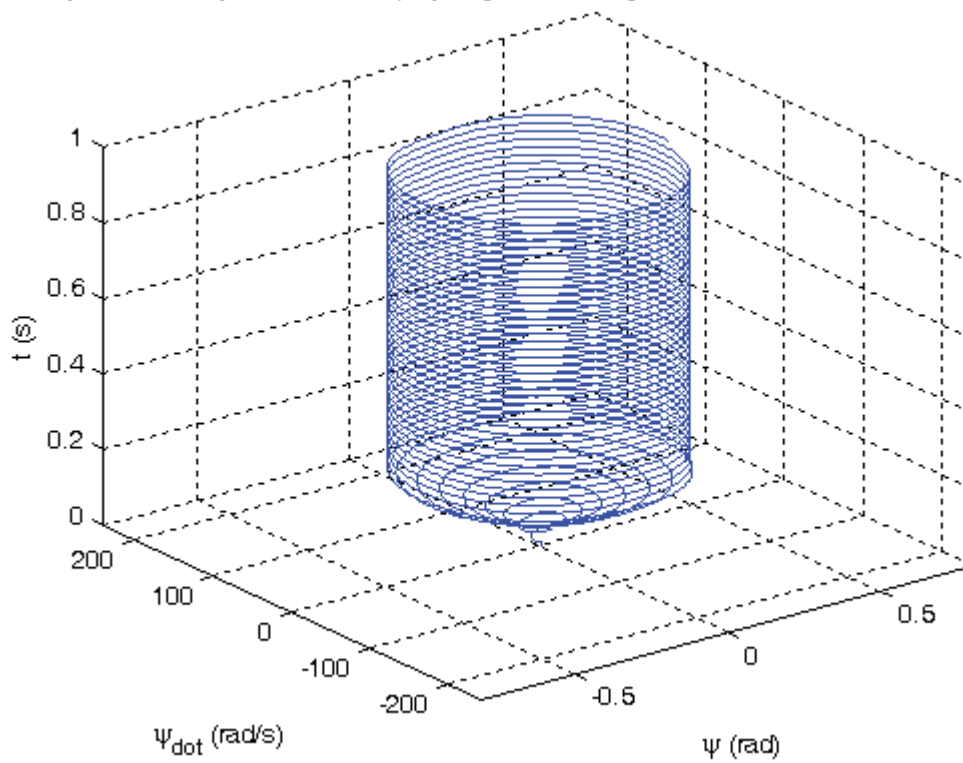
Limit cycle for the system with freeplay angle of 1 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



1718

Figure 22: Figure 17 :Figure 18 :

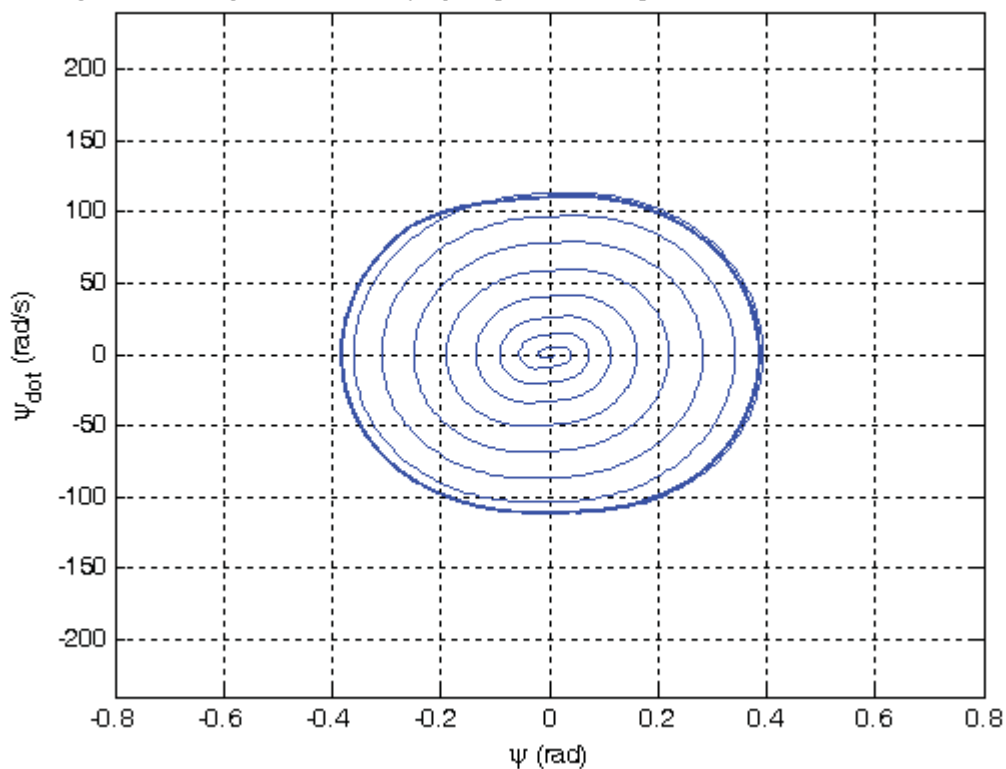
Limit cycle for the system with freeplay angle of 1.5 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



19

Figure 23: Figure 19 :

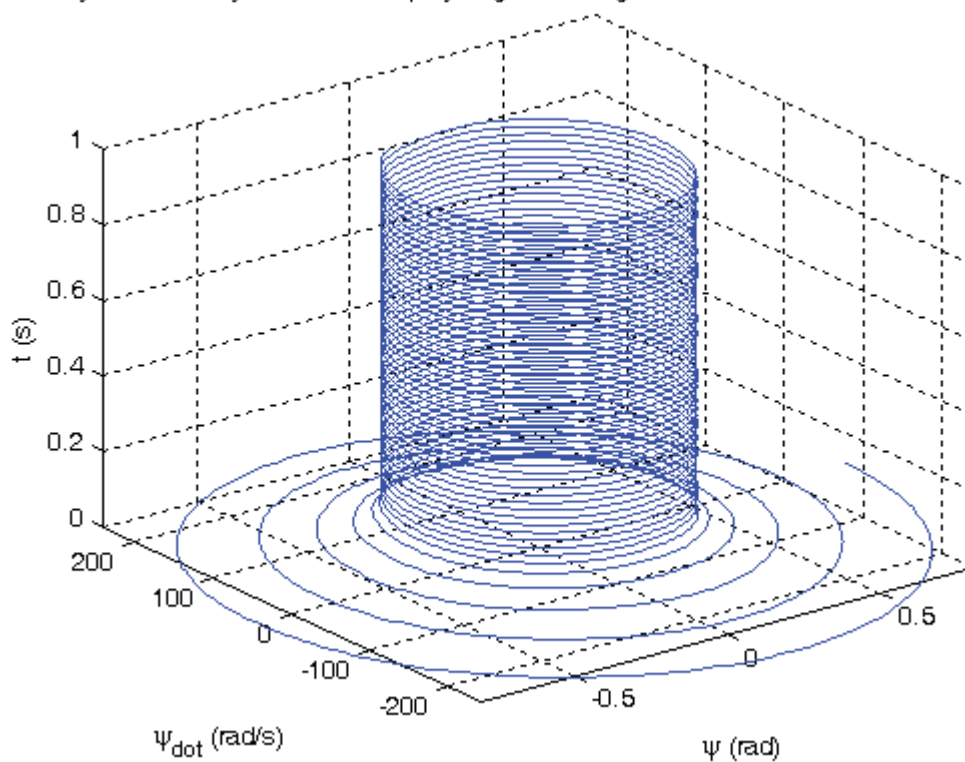
Limit cycle for the system with freeplay angle of 1.5 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



20

Figure 24: Figure 20 :

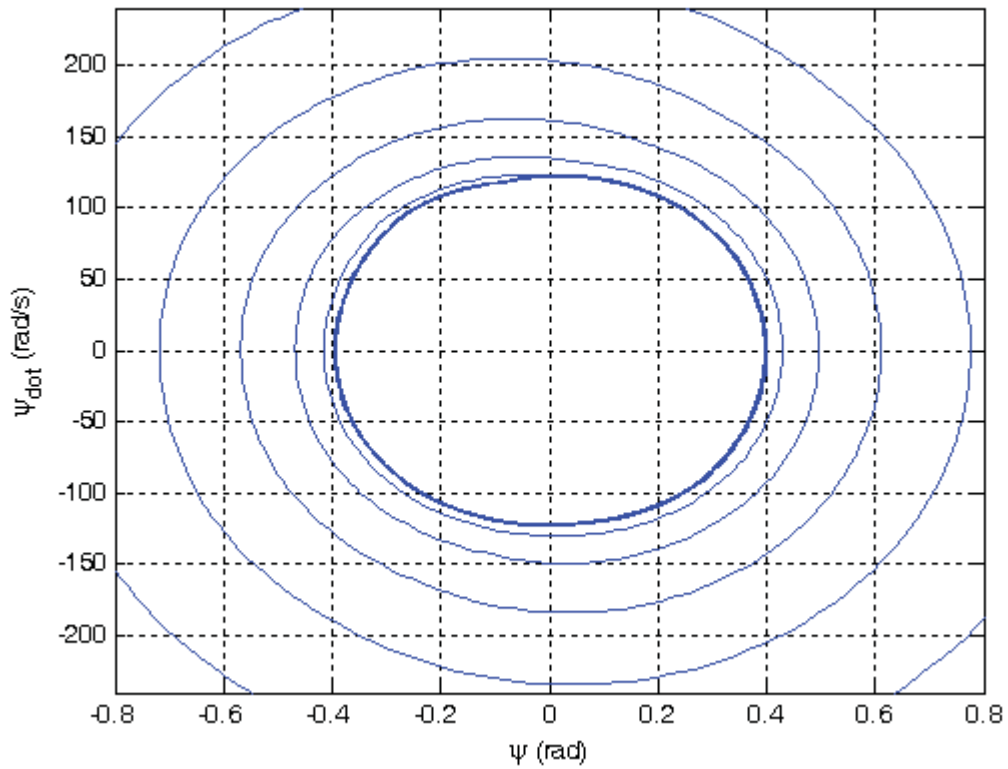
Limit cycle for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



21

Figure 25: Figure 21 :

Limit cycle for the system with freeplay angle of 0 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s



22

Figure 26: Figure 22 :

Limit cycle for the system with freeplay angle of 0.5 degrees for  $v=50$  m/s and  $k=-10$  Nm/rad/s

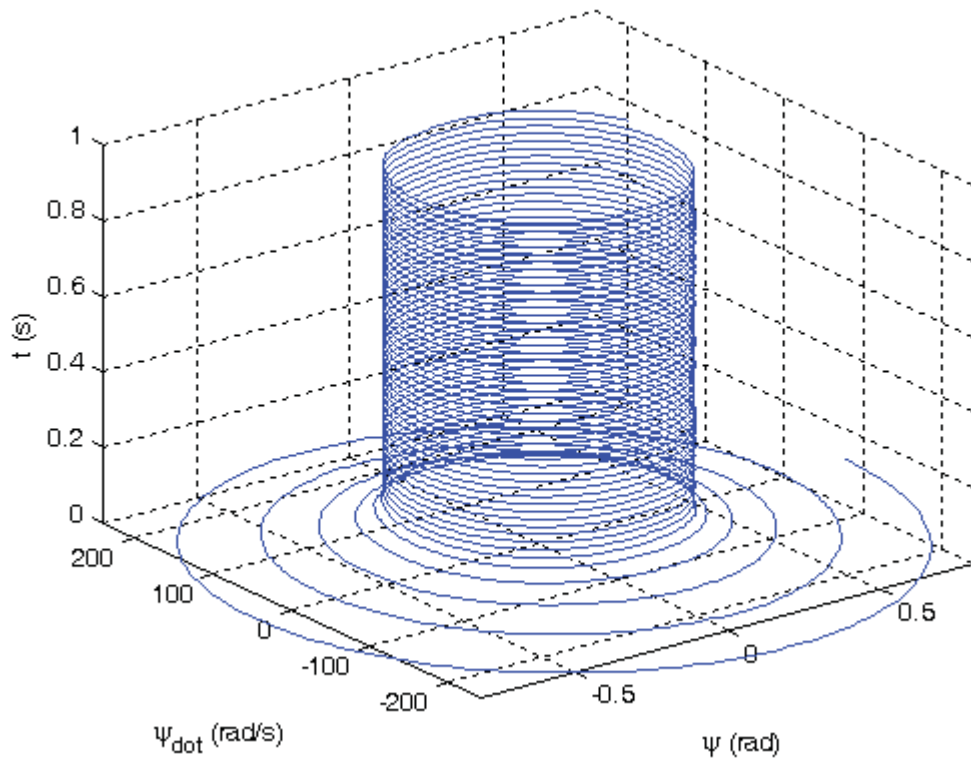


Figure 27: Volume

*[Note: 30. Chartier, B., Tuohy, B., Retallack, J., Tennant, S., -: Landing gear shock absorber. Research project. [ftp://ftp.uniduisburg.de/FlightGear/Docs/Landing\\_Gear\\_Shock\\_Absorber.pdf](ftp://ftp.uniduisburg.de/FlightGear/Docs/Landing_Gear_Shock_Absorber.pdf) accessed on November 23]*

Figure 28:

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